

Une preuve constructive d'existence d'orbites périodiques spontanées pour les équations de Navier-Stokes

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Spontaneous periodic orbits of the Navier-Stokes equations

- ▶ Consider the Navier-Stokes equations on the torus

$$\begin{cases} \partial_t u + (u \cdot \nabla)u - \nu \Delta u + \nabla p = f & \text{on } \mathbb{R} \times \mathbb{T}^3 \\ \nabla \cdot u = 0 & \text{on } \mathbb{R} \times \mathbb{T}^3 \end{cases}$$

with a *Taylor-Green* type of forcing

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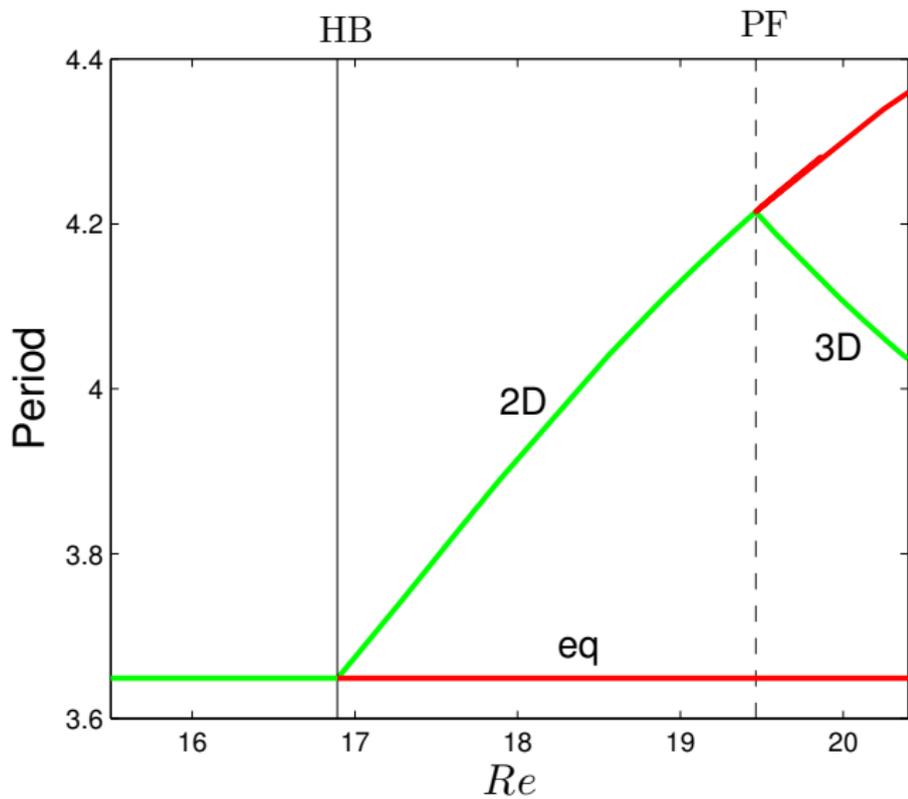
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- ▶ This steady state is stable if the fluid is viscous enough. When ν decreases, it becomes unstable and the dynamics becomes more and more complex.



A few references

- ▶ In the presence of a periodic external influence, periodic motions in fluids have been studied extensively, and are relatively well understood [Serrin '59; Kaniel & Shinbrot '67; Takeshita '69; Maremonti '91; Kozono & Nakao '96; Yamazaki '00; Galdi & Sohr '04; etc.].

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- ▶ In the absence of a periodic external influence, periodic motions in fluids are much harder to study, and the existing results are typically of perturbative nature [Iudovich '71; Iooss '72; Joseph & Sattinger '72; Melcher, Schneider & Uecker '08; Galdi '16].

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Can we say anything rigorous about this specific “solution”?

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- ▶ We want an *a posteriori* error bound, but without knowing *a priori* that the true zero exists.

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A Newton-Kantorovich type of theorem

Let \mathcal{X}, \mathcal{Y} be Banach spaces, $F : \mathcal{X} \rightarrow \mathcal{Y}$ a \mathcal{C}^1 function. Let $\bar{x} \in \mathcal{X}$ and assume we have the following estimates:

$$\begin{aligned}\|F(\bar{x})\|_{\mathcal{Y}} &\leq \varepsilon \\ \|DF(\bar{x})^{-1}\|_{\mathcal{L}(\mathcal{Y}, \mathcal{X})} &\leq \kappa \\ \|DF(x) - DF(\bar{x})\|_{\mathcal{L}(\mathcal{X}, \mathcal{Y})} &\leq \gamma (\|x - \bar{x}\|_{\mathcal{X}}).\end{aligned}$$

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If there exists $r > 0$ such that

$$\kappa\varepsilon + \kappa\gamma(r)r < r,$$

then F has a unique zero x satisfying $\|x - \bar{x}\|_{\mathcal{X}} \leq r$.

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Proof.

The operator $T : x \mapsto x - DF(\bar{x})^{-1}F(x)$ is a contraction on $B(\bar{x}, r)$. \square

A Newton-Kantorovich type of theorem, version 2

Let \mathcal{X}, \mathcal{Y} be Banach spaces, $F : \mathcal{X} \rightarrow \mathcal{Y}$ a \mathcal{C}^1 function. Let $\bar{x} \in \mathcal{X}$, $A \in \mathcal{L}(\mathcal{Y}, \mathcal{X})$ injective, and assume we have the following estimates:

$$\|F(\bar{x})\|_{\mathcal{Y}} \leq \varepsilon$$

$$\|A\|_{\mathcal{L}(\mathcal{Y}, \mathcal{X})} \leq \kappa$$

$$\|I - ADF(\bar{x})\|_{\mathcal{L}(\mathcal{X}, \mathcal{X})} \leq \delta < 1$$

$$\|DF(x) - DF(\bar{x})\|_{\mathcal{L}(\mathcal{X}, \mathcal{Y})} \leq \gamma (\|x - \bar{x}\|_{\mathcal{X}}).$$

If there exists $r > 0$ such that

$$\kappa\varepsilon + (\delta + \kappa\gamma(r))r < r,$$

then F has a unique zero x satisfying $\|x - \bar{x}\|_{\mathcal{X}} \leq r$.

Proof.

The operator $T : x \mapsto x - AF(x)$ is a contraction on $B(\bar{x}, r)$. □

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- ▶ We want to use this procedure to study periodic orbits of the Navier-Stokes equations.
- ▶ The first step is to find an appropriate $F = 0$ formulation, on a well-chosen Banach space.

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- We consider $\omega = \nabla \times u$, $g = \nabla \times f$, and apply the curl to NS

$$\begin{cases} \partial_t \omega + (u \cdot \nabla)\omega - (\omega \cdot \nabla)u - \nu \Delta \omega = g \\ \nabla \cdot u = 0 \end{cases}$$

Reformulation using the vorticity

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$$\begin{cases} \partial_t \omega + (u \cdot \nabla)\omega - (\omega \cdot \nabla)u - \nu \Delta \omega = g \\ \nabla \cdot u = 0 \end{cases}$$

- ▶ Thanks to the continuity equation, we can express u as a function of ω

$$u = -\Delta^{-1} \nabla \times \omega,$$

which yields

$$\partial_t \omega - ((\Delta^{-1} \nabla \times \omega) \cdot \nabla)\omega + (\omega \cdot \nabla)(\Delta^{-1} \nabla \times \omega) - \nu \Delta \omega = g.$$

Looking for periodic solutions

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- ▶ We call Ω the (unknown) frequency of the solution, and write

$$\omega(t, x) = \sum_{n=(n_t, \tilde{n}) \in \mathbb{Z}^4} \omega_n e^{i(n_t \Omega t + \tilde{n} \cdot x)}$$

with $\tilde{n} = (n_x, n_y, n_z)$ the space indices and n_t the time index.

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- ▶ Denoting $W = (\Omega, (\omega_n)_{n \in \mathbb{Z}^4})$ and Fourier-transforming the vorticity equation, we obtain the problem $F(W) = (F_n(W))_{n \in \mathbb{Z}^4} = 0$ with

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 - Hence, $F_n(\bar{W})$ is nonzero for only finitely many n , and $\|F(\bar{W})\|$ can be computed explicitly (using interval arithmetic)
- The computation of γ satisfying

$$\left\| DF(W) - DF(\bar{W}) \right\| \leq \gamma \left(\|W - \bar{W}\| \right)$$

is rather straightforward (Banach algebra property).

The key estimate : $\|I - ADF(\bar{W})\| < 1$

- ▶ The asymptotically dominant terms in $DF(\bar{W})$ are given by the eigenvalues of the heat operator:

$$\lambda_n = i\bar{\Omega}n_t + \nu(n_x^2 + n_y^2 + n_z^2).$$

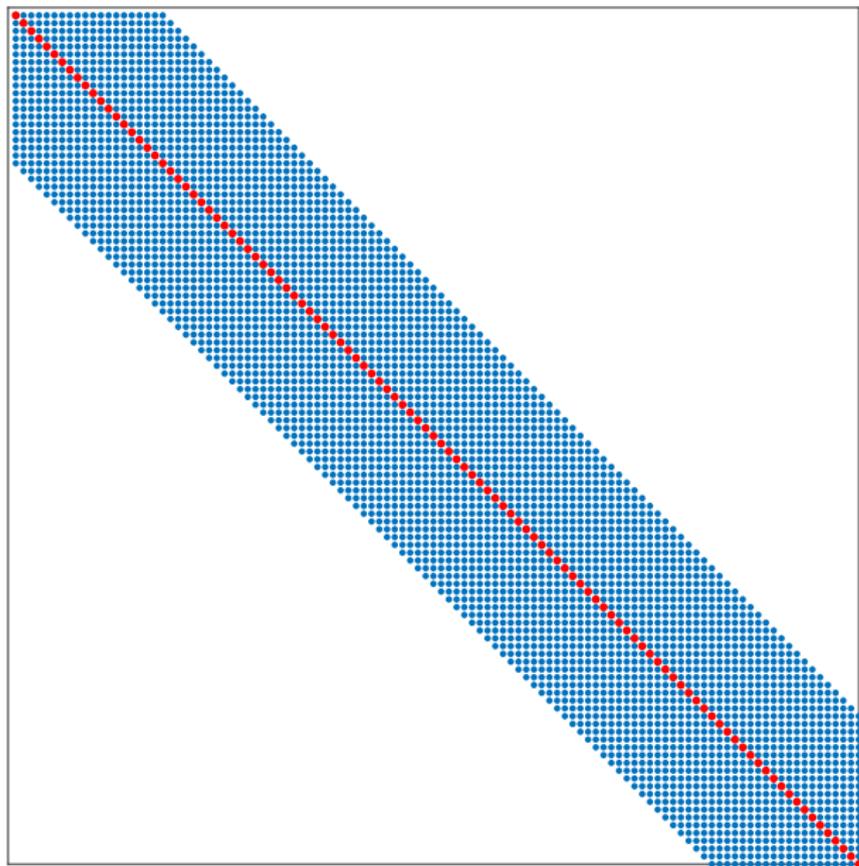
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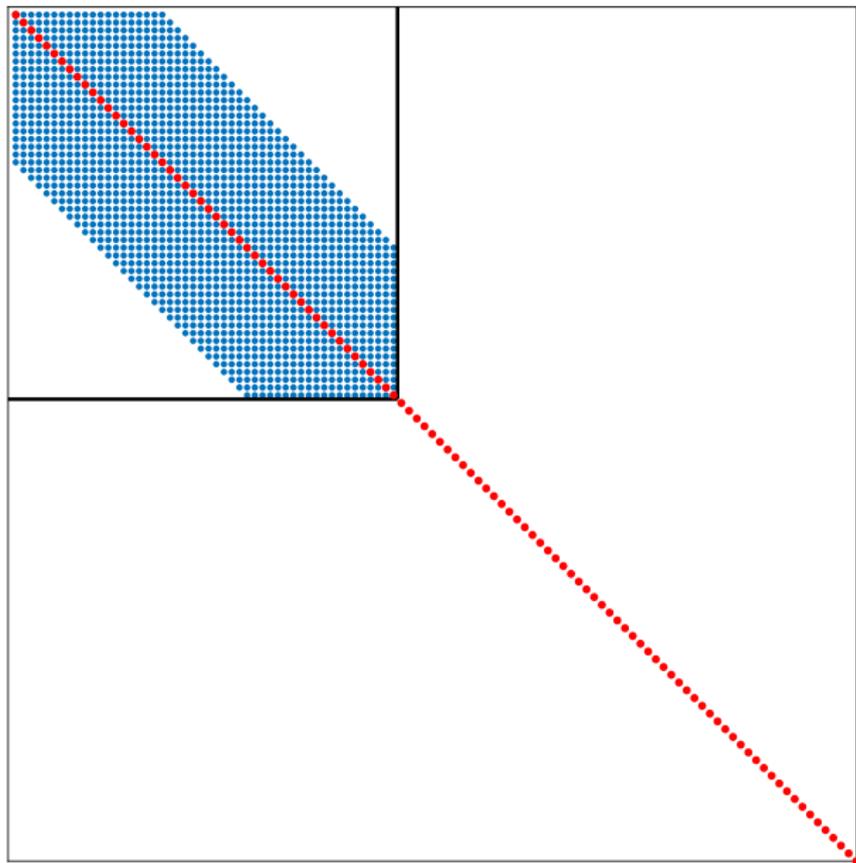
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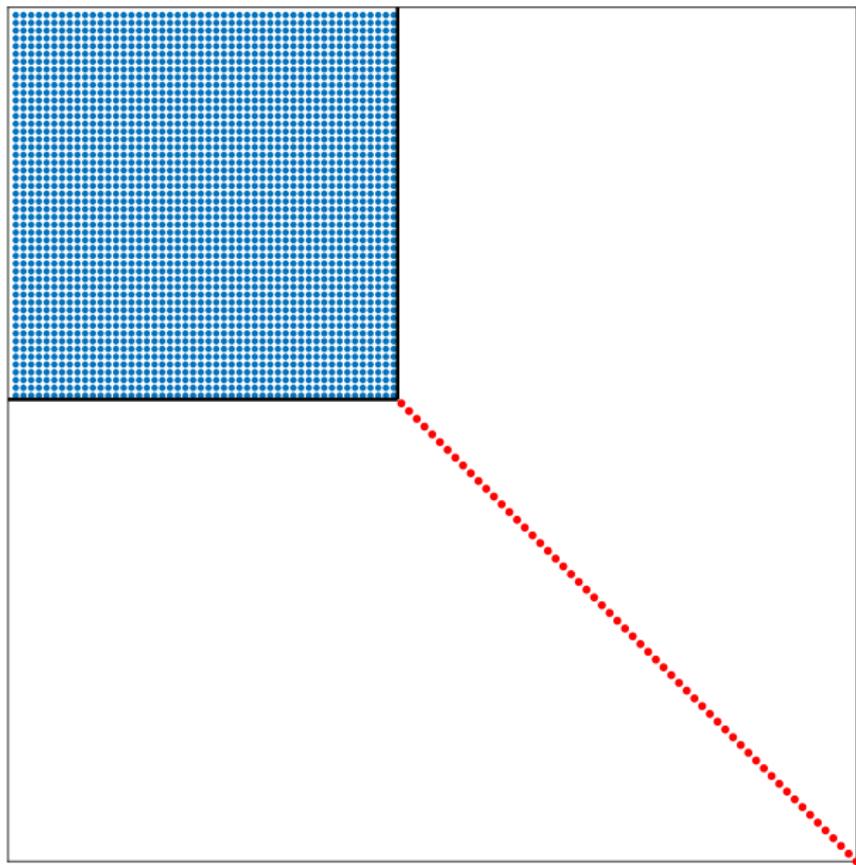
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The inverse eigenvalues: $\lambda_n^{-1} = \left(i\bar{\Omega}n_t + \nu(n_x^2 + n_y^2 + n_z^2) \right)^{-1}$

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- ▶ The finite dimensional part can be obtained by inverting some finite dimensional projection of $DF(\bar{x})$.

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- ▶ In practice, it is crucial to reduce as much as possible the dimension of the finite-dimensional space that we keep for the validation (which drastically increases with m).

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- ▶ These symmetries can first be used to reduce the number of Fourier modes used for the numerical simulations.
- ▶ By deriving a posteriori estimates that are *compatible* with the symmetries, we can also reduce the number of modes used for the validation (and in particular to reduce the size of the finite block used in A).
- ▶ It turns out that the first branch of periodic orbits that we obtain after the bifurcation do not depend on z (the solutions are essentially 2D), which we also use to reduce the number of modes.

Contour lines of the vertical vorticity $\omega^{(z)}$. $\nu = 0.286$.

Theorem [B., Lessard, van den Berg & van Veen '21]

There exists a periodic solution of NS at a distance of at most 10^{-5} (in C^0 norm) of this numerical solution.

THANK YOU FOR YOUR ATTENTION!