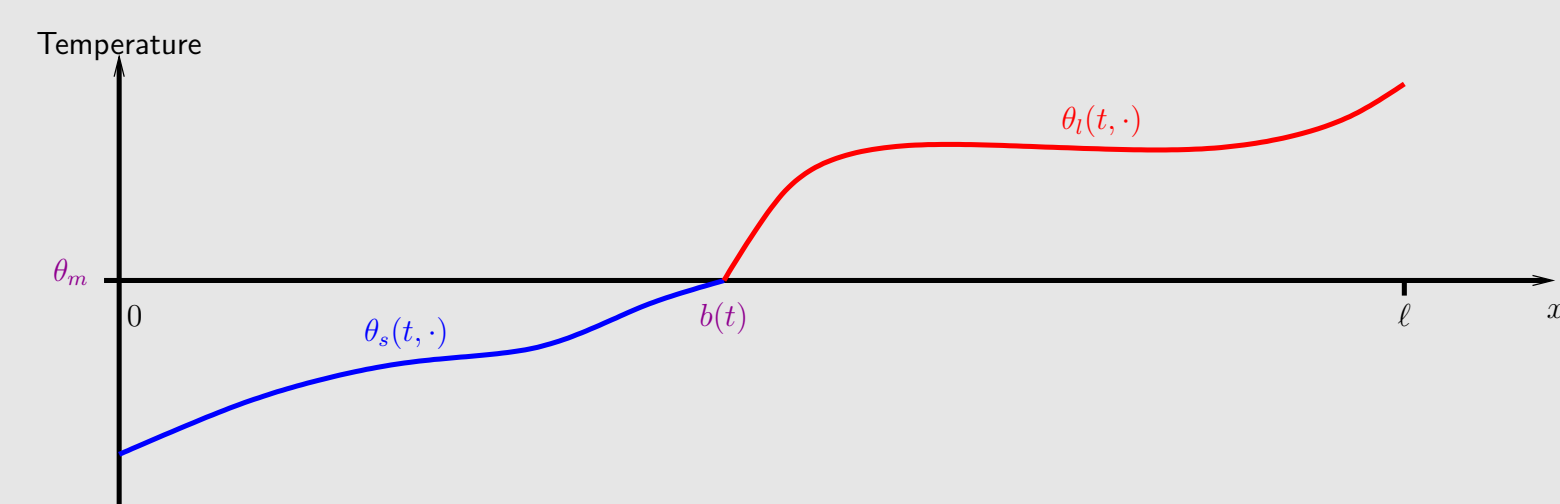
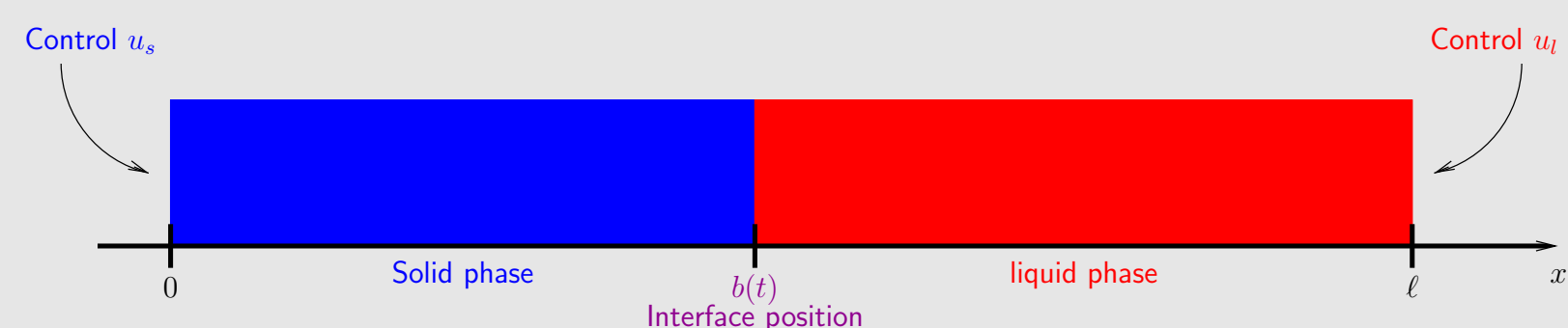


Introduction

In this poster, we are interested in the controllability of the **Stefan problem**, which is a mathematical model for **solid-liquid phase transition**, well-posedness of the problem has already been studied intensively, see *Fasano and Primicerio 1977* for example. Here, we present a **steady state to steady state controllability result** based on the **flatness method** which has been applied to the **heat equation** in *Martin et al. 2014* and to the **Stefan problem** in *Dunbar et al. 2003*.



Problem statement

If we denote by θ_s and θ_l **the temperature in the two slabs** (solid and liquid) and by b **the position of the moving interface between the two**, using dimensionless variables, the **Stefan problem** can be written as:

$$\dot{\theta}_s(t, x) = \partial_x^2 \theta_s(t, x) \quad (t > 0, x \in (0, b(t))), \quad (1a)$$

$$\theta_s(t, b(t)) = 0 \quad (t > 0), \quad (1b)$$

$$\theta_s(t, 0) = u_s(t) \quad (t > 0), \quad (1c)$$

$$\dot{\theta}_l(t, x) = c \partial_x^2 \theta_l(t, x) \quad (t > 0, x \in (b(t), 1)), \quad (2a)$$

$$\theta_l(t, b(t)) = 0 \quad (t > 0), \quad (2b)$$

$$\theta_l(t, 1) = u_l(t) \quad (t > 0), \quad (2c)$$

$$\dot{b}(t) = \partial_x \theta_s(t, b(t)) - \partial_x \theta_l(t, b(t)) \quad (t > 0), \quad (3)$$

with some **initial conditions**,

$$b(0) = b^0, \quad (4a)$$

$$\theta_s(0, x) = \theta_s^0(x) \quad (x \in (0, b^0)), \quad (4b)$$

$$\theta_l(0, x) = \theta_l^0(x) \quad (x \in (b^0, 1)), \quad (4c)$$

with the **state constraints**,

$$b(t) \in [0, 1] \quad (t > 0), \quad (5a)$$

$$\theta_s(t, x) \leq 0 \quad (t > 0, x \in (0, b(t))), \quad (5b)$$

$$\theta_l(t, x) \geq 0 \quad (t > 0, x \in (b(t), 1)), \quad (5c)$$

and with the **control constraints**,

$$u_s(t) \leq 0 \quad \text{and} \quad u_l(t) \geq 0 \quad (t > 0). \quad (6)$$

Let us mention that **the steady states** of this system are described by $\bar{b} \in (0, 1)$ and a parameter $\bar{v} \in \mathbb{R}_+$, and given $\bar{b} \in (0, 1)$ and $\bar{v} \in \mathbb{R}_+$, the associated **steady state** is given by

$$\bar{u}_s = -\bar{v}\bar{b}, \quad \bar{u}_l = \bar{v}(1 - \bar{b}),$$

$$\bar{\theta}_s(x) = \bar{v}(x - \bar{b}) \quad (x \in [0, \bar{b}]) \quad \text{and} \quad \bar{\theta}_l(x) = \bar{v}(x - \bar{b}) \quad (x \in [\bar{b}, 1]). \quad (7)$$

Let us then define the set of **steady states**:

$$\mathcal{S}_+^*(\bar{b}) = \left\{ (\bar{\theta}_s, \bar{\theta}_l) \in \left(H^2(0, \bar{b}) \cap H_R^1(0, \bar{b}) \right) \times \left(H^2(\bar{b}, 1) \cap H_L^1(\bar{b}, 1) \right) \mid \right.$$

$$\left. \exists \bar{v} \in \mathbb{R}_+^*, \forall x \in (0, \bar{b}), \theta_s(x) = \bar{v}(x - \bar{b}) \text{ and } \forall x \in (\bar{b}, 1), \theta_l(x) = \bar{v}(x - \bar{b}) \right\}. \quad (8a)$$

Main result: Steady state to steady state controllability theorem

Let $\bar{b}_0, \bar{b}_1 \in (0, 1)$, $(\bar{\theta}_s^0, \bar{\theta}_l^0)$ and $(\bar{\theta}_s^1, \bar{\theta}_l^1)$ be steady states, and let $\bar{v}_0, \bar{v}_1 \in \mathbb{R}_+$ be such that

$$\bar{\theta}_s^0(x) = \bar{v}_0(x - \bar{b}_0) \quad (x \in [0, \bar{b}_0]) \quad \text{and} \quad \bar{\theta}_l^0(x) = \bar{v}_0(x - \bar{b}_0) \quad (x \in [\bar{b}_0, 1])$$

and

$$\bar{\theta}_s^1(x) = \bar{v}_1(x - \bar{b}_1) \quad (x \in [0, \bar{b}_1]) \quad \text{and} \quad \bar{\theta}_l^1(x) = \bar{v}_1(x - \bar{b}_1) \quad (x \in [\bar{b}_1, 1]).$$

Then for all $T > 0$, there exists $u_s, u_l \in C^\infty(0, T)$, Gevrey, such that the solution of (1)–(3) is steered from $(\bar{\theta}_s^0, \bar{\theta}_l^0, \bar{b}^0)$ to $(\bar{\theta}_s^1, \bar{\theta}_l^1, \bar{b}^1)$ in time T .

In addition, if $(\bar{\theta}_s^0, \bar{\theta}_l^0) \in \mathcal{S}_+^*(\bar{b}^0)$ and $(\bar{\theta}_s^1, \bar{\theta}_l^1) \in \mathcal{S}_+^*(\bar{b}^1)$ (i.e., if $\bar{v}_0, \bar{v}_1 \in \mathbb{R}_+^*$), then there exists $T > 0$ and $u_s, u_l \in C^\infty(0, T)$, satisfying (5)–(6) such that the solution of (1)–(3) is steered from $(\bar{\theta}_s^0, \bar{\theta}_l^0, \bar{b}^0)$ to $(\bar{\theta}_s^1, \bar{\theta}_l^1, \bar{b}^1)$ in time T .

Gevrey functions

Let $n \geq 1$, $U \subset \mathbb{R}^n$ and $f \in C^\infty(U, \mathbb{R})$, we say that f is **Gevrey**, if there exist $M \geq 0$, $R_1, \dots, R_n > 0$ and $\sigma_1, \dots, \sigma_n \geq 0$ such that:

$$\forall x \in U, \forall p_1, \dots, p_n \in \mathbb{N}, \quad \left| \partial_{x_1}^{p_1} \dots \partial_{x_n}^{p_n} f(x) \right| \leq M \prod_{i=1}^n \frac{(p_i)!^{\sigma_i}}{R_i^{p_i}}$$

And if $p \in \mathbb{N}$ and $f \in C^\infty(U, \mathbb{R}^p)$, we say that f is **Gevrey** if its coordinates are **Gevrey** in the above sense. If we are working on an interval I of \mathbb{R} , we denote $\mathcal{G}(M, R, \sigma)$ the set of Gevrey functions on I of order σ and constants $M, R \geq 0$

Estimate Result

Let $\chi \geq 0$, $I \subset \mathbb{R}$ an interval, $\sigma \in [1, 2]$, $M_\alpha > 0$, $M_\beta, M_f \geq 0$ and $R > 0$. Let $f \in \mathcal{G}(M_f, R, \sigma)$, $\alpha_0 \in \mathcal{G}(M_\alpha, R, \sigma)$ and $\beta_0 \in \mathcal{G}(M_\beta, R, \sigma)$ be given Gevrey functions defined on I . Consider the sequence defined by

$$\begin{cases} \beta_{i+1} = \chi \dot{\beta}_i - \chi f \alpha_i, \\ \alpha_{i+1} = \chi \dot{\alpha}_i - \chi f \beta_{i+1}, \end{cases}$$

where β_i and α_i are real functions defined on I (we initialize the sequences with α_0, β_0).

Then, for every $i \in \mathbb{N}$, α_i and β_i are Gevrey functions of order σ defined on I . In addition, for every $\rho \in (0, \rho^*(M, R, \chi)]$, we have, for every $l \in \mathbb{N}$ and every $i \in \mathbb{N}$,

$$\|\alpha_i^{(l)}\|_{L^\infty(I)} \leq \frac{M_\alpha \chi^i (l + 2i)!^\sigma}{R^l \rho^i i!^{p(2i)!^{\sigma-1}}} \quad \text{and} \quad \|\beta_{i+1}^{(l)}\|_{L^\infty(I)} \leq \frac{\mu \chi^{i+1} (l + 2i + 1)!^\sigma}{R^l \rho^i i!^{p(2i + 1)!^{\sigma-1}}}, \quad (9)$$

where we have set $\mu := \frac{M_\beta}{R} + M_\alpha M_f$, $p = 2 - \sigma \in [0, 1]$ and

$$\rho^*(M, R, \chi) = \min \left\{ \left(\frac{4 + \chi M_\beta M_f / (2M_\alpha)}{R} + \frac{\chi M_f^2}{2} \right)^{-1}, \frac{2R}{3} \right\}.$$

Step and bump functions

If $\sigma > 1$ and $k := (\sigma - 1)^{-1}$ the function defined by:

$$\phi_\sigma(t) := \begin{cases} 1 & \text{if } t \leq 0, \\ 0 & \text{if } t \geq 1, \\ \frac{e^{-(1-t)^{-k}}}{e^{-(1-t)^{-k}} + e^{-t^{-k}}} & \text{if } t \in (0, 1), \end{cases} \quad (t \in \mathbb{R}), \quad (10)$$

is **Gevrey** of order σ on \mathbb{R} and verifies:

$$\mathbf{1} \quad \forall i \in \mathbb{N}^*, \phi_\sigma^{(i)}(0) = \phi_\sigma^{(i)}(1) = 0$$

$$\mathbf{2} \quad \forall t \in \mathbb{R}, 0 \leq \phi_\sigma(t) \leq 1.$$

$$\mathbf{3} \quad \forall t \in \mathbb{R}, \phi_\sigma(t) + \phi_\sigma(1 - t) = 1$$

If $\sigma > 1$ then, it exists η a **Gevrey** function of order σ on \mathbb{R} such that:

$$\mathbf{1} \quad \eta \text{ vanishes on } \mathbb{R} \setminus [0, 1]$$

$$\mathbf{2} \quad \eta \geq 0$$

$$\mathbf{3} \quad \int_{\mathbb{R}} \eta(t) dt = 1$$

It is constructed using the step function.

Sketch of the proof of the main result

We are looking for solutions under the form:

$$\theta_s(t, x) = \sum_{i=0}^{\infty} \alpha_i^s(t) \frac{(x - b(t))^{2i+1}}{(2i+1)!} + \sum_{i=0}^{\infty} \beta_i^s(t) \partial_x \frac{(x - b(t))^{2i}}{(2i)!}, \quad (11a)$$

$$\theta_l(t, x) = \sum_{i=1}^{\infty} \alpha_i^l(t) \frac{(x - b(t))^{2i+1}}{(2i+1)!} + \sum_{i=0}^{\infty} \beta_i^l(t) \partial_x \frac{(x - b(t))^{2i}}{(2i)!}. \quad (11b)$$

The problem statement gives us the following conditions:

$$\begin{cases} \beta_{i+1}^s = \dot{\beta}_i^s - \dot{b} \alpha_i^s, \\ \alpha_{i+1}^s = \dot{\alpha}_i^s - \dot{b} \beta_{i+1}^s \end{cases} \quad \text{and} \quad \begin{cases} c \beta_{i+1}^l = \dot{\beta}_i^l - \dot{b} \alpha_i^l, \\ c \alpha_{i+1}^l = \dot{\alpha}_i^l - \dot{b} \beta_{i+1}^l \end{cases} \quad (12)$$

and

$$\dot{b} = \alpha_0^s - \alpha_0^l. \quad (13)$$

The previous estimate result tells us that if α_0^s and α_0^l are chosen in some $\mathcal{G}(M, R, \sigma)$ with $\sigma \in (1, 2)$ then the previous series converge and are \mathcal{C}^∞ and are in fact Gevrey (we necessarily have $\beta_0^s = \beta_0^l = 0$).

If $b_1 \geq b_0$ (else we switch the two), we set:

$$\alpha_0^l(t) := \alpha\left(\frac{t}{T}\right) \quad \text{and} \quad \alpha_0^s(t) := \alpha\left(\frac{t}{T}\right) + \frac{1}{T} \varphi\left(\frac{t}{T}\right) \quad (t \in [0, T]). \quad (14)$$

with

$$\alpha(t) = \bar{v}_0 \phi_\sigma(t) + \bar{v}_1 \phi_\sigma(1 - t) \quad \text{and} \quad \varphi(t) = (\bar{b}_1 - \bar{b}_0) \eta(t) \quad (t \in [0, 1]),$$

to have

$$\mathbf{■} \quad \alpha_0^s(0) = \alpha_0^l(0) = \bar{v}_0;$$

$$\mathbf{■} \quad \alpha_0^s(T) = \alpha_0^l(T) = \bar{v}_1;$$

$$\mathbf{■} \quad \alpha_0^{s(i)}(0) = \alpha_0^{l(i)}(0) = \alpha_0^{s(i)}(T) = \alpha_0^{l(i)}(T) = 0 \text{ for every } i \in \mathbb{N};$$

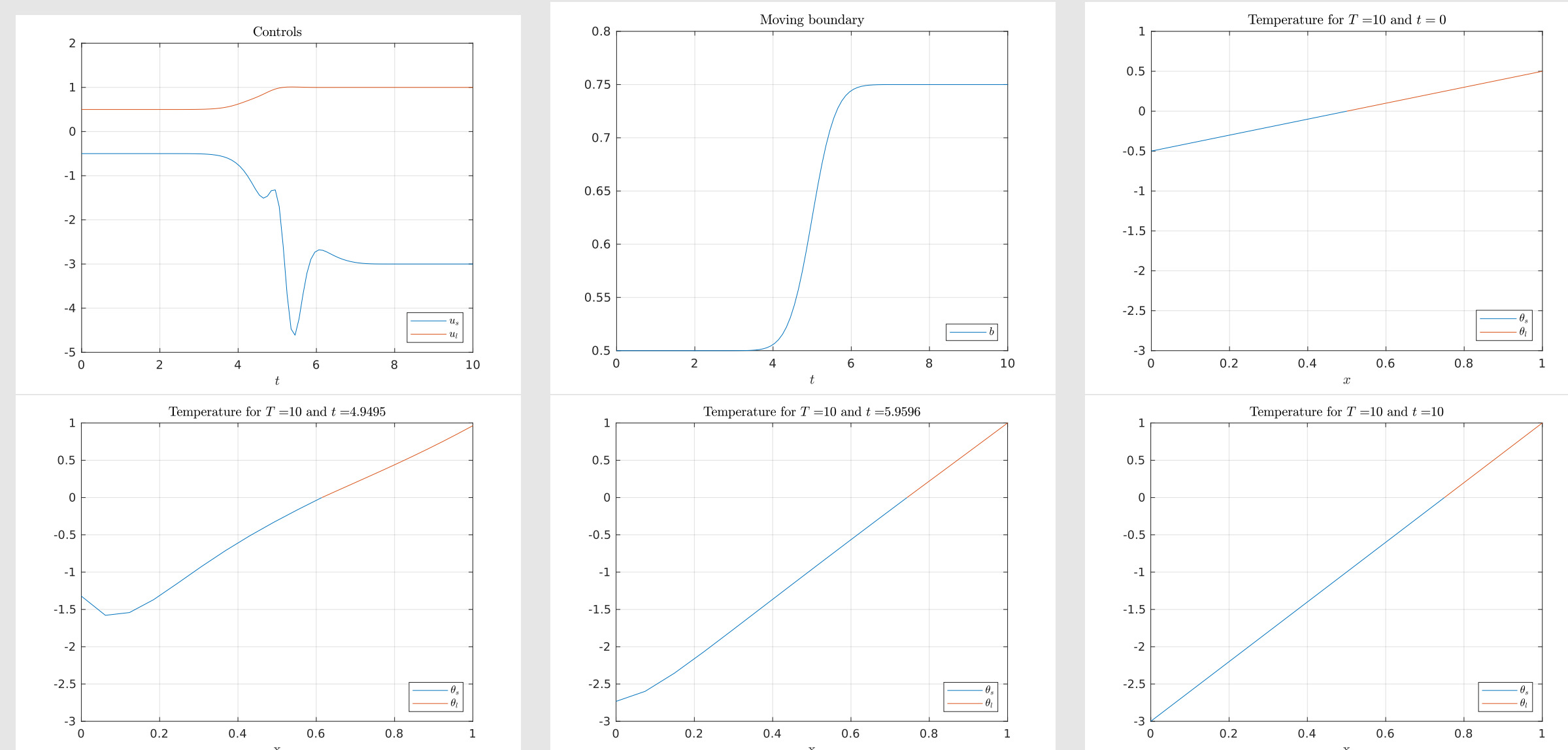
$$\mathbf{■} \quad \int_0^T (\alpha_0^s(t) - \alpha_0^l(t)) dt = \bar{b}_1 - \bar{b}_0.$$

And thus we designed a solution to our problem. We get u_s and u_l as the traces of our solutions.

To show the preservation of the sign constraints for large T we derive finer estimates on our series depending on T , using the construction of α_0^s and α_0^l and our estimation result.

Numerical aspect

The series defining the trajectories we designed in our result converge at a rate: $o\left(\frac{e^{-\frac{2}{3}N(\ln(N)-1)}}{(N+1)^2 \ln(N)}\right)$.



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