

# Méthode de Newton pour le calcul de carènes optimales basée sur la formule de Michell pour des vitesses aléatoires

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CANUM

1. The problem
2. Expectation of the water resistance
3. Newton's method
4. Numerical results

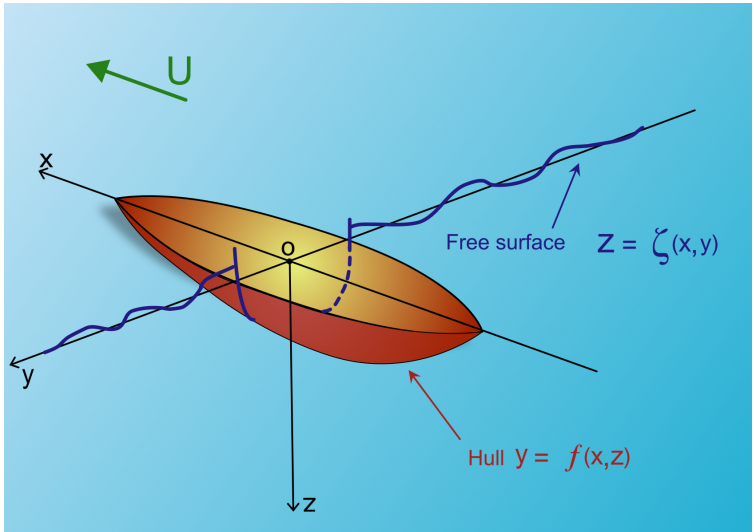
# The setting

## The problem

Expectation of the  
water resistance

Newton's method

Numerical results



## The problem

Expectation of the  
water resistance

Newton's method

Numerical results

The fluid is assumed to be

- 1 **incompressible, inviscid** : We start from the Euler equations :

$$\begin{cases} \partial_t \mathbf{v} + (\mathbf{v} \cdot \nabla) \mathbf{v} = \frac{1}{\rho} \nabla P + \mathbf{g} & \text{in } W = (\text{Domain of the water}) \\ \nabla \cdot \mathbf{v} = 0 \end{cases}$$

- 2 **irrotational** :

- $\nabla \times \mathbf{v} = 0$  in  $W$  .
- $\exists \Phi, \mathbf{v} = \nabla \Phi$  (Helmholtz-Hodge Theorem).

- 3 A steady state has been reached so that we have

$$\Delta \Phi = 0 \quad \text{in } W.$$



# The problem

## The problem

### Expectation of the water resistance

### Newton's method

### Numerical results

If the ship hull is represented by  $y = f(x, z)$ , and  $z = \zeta(x, y)$  is the unknown free surface of the water, the boundary conditions are :

- Hull boundary condition :

$$\nabla \Phi \cdot n_f = 0 \quad \text{on } y = f(x, z)$$

- Free surface conditions :

- i Kinematic condition :

$$\nabla \Phi \cdot n_\zeta = 0 \quad \text{on } z = \zeta(x, y),$$

- ii Dynamic condition (Bernoulli's equation) :

$$|\nabla \Phi|^2 - 2g\zeta + \frac{2}{\rho}P = \frac{2}{\rho}P_0 \quad \text{on } z = \zeta(x, y)$$

## The problem

## Expectation of the water resistance

## Newton's method

## Numerical results

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$$\begin{cases} \Delta \phi &= 0, \\ \partial_{xx} \phi - \frac{g}{U^2} \partial_z \phi &= 0, \\ \partial_y \phi - U \partial_x f &= 0, \end{cases} \quad \begin{array}{l} \text{in } \mathbb{R} \times \mathbb{R}^+ \times \mathbb{R}^+ \quad (\text{quarter space}) \\ \text{on } z = 0 \\ \text{on } y = 0^+ \end{array}$$

- $\Phi = -Ux + \phi \quad (\phi \text{ small}) .$

- Thin-ship assumptions :

$$|\partial_x f| \ll 1 \quad ; \quad |\partial_z f| \ll 1$$

- Michell condition on the hull :

$$\partial_y \phi - U \partial_x f(x, z) = 0 \quad \text{on } y = 0^+$$

- Linearized free surface conditions :

- i Kinematic condition :

$$\partial_z \phi + U \partial_x \zeta = 0 \quad \text{on } z = 0$$

- ii Dynamic condition :

$$\partial_{xx} \phi + \frac{g}{U} \partial_x \zeta = 0 \quad \text{on } z = 0$$

# Solution and wave resistance

- If  $\delta p$  is the increase of fluid pressure due to the disturbance  $\Phi$ , by Bernoulli's equation, the wave resistance is given by

$$\begin{aligned} R_{\text{Michell}} &= -2 \iint \delta p \partial_x f dx dz = 2\rho U \iint \partial_x \phi \partial_x f dx dz, \\ &= \frac{4\rho g^2}{\pi U^2} \int_1^\infty (\mathcal{I}^2(\lambda) + \mathcal{J}^2(\lambda)) \frac{\lambda^2 d\lambda}{\sqrt{\lambda^2 - 1}} \end{aligned}$$

where, if  $D$  denotes the domain of definition of  $f$ , we have :

$$\begin{aligned} \mathcal{I}(\lambda) &= \int_D \partial_x f(x, z) \exp(-\lambda^2 gz/U^2) \cos(\lambda gx/U^2) dx dz, \\ \mathcal{J}(\lambda) &= \int_D \partial_x f(x, z) \exp(-\lambda^2 gz/U^2) \sin(\lambda gx/U^2) dx dz. \end{aligned}$$

- A standard approximation of the viscous drag for small  $\nabla f$  is given by a linearization of the area functional and reads

$$R_{\text{viscous}} = \frac{1}{2} \rho U^2 C_F \left( 2|D| + \int_D |\nabla f(x, z)|^2 dx dz \right)$$

The problem

Expectation of the  
water resistance

Newton's method

Numerical results

## The problem

## Expectation of the water resistance

## Newton's method

## Numerical results

The normalized total water resistance writes :  $J_{total} = J_{Michell} + J_{viscous}$ ,  
with

$$J_{viscous} = \int_D |\nabla f(x, z)|^2 dx dz$$

$$J_{Michell} = \int_{D \times D} k(U, x, z, x', z') f(x, z) f(x', z') dx dz dx' dz'$$

where

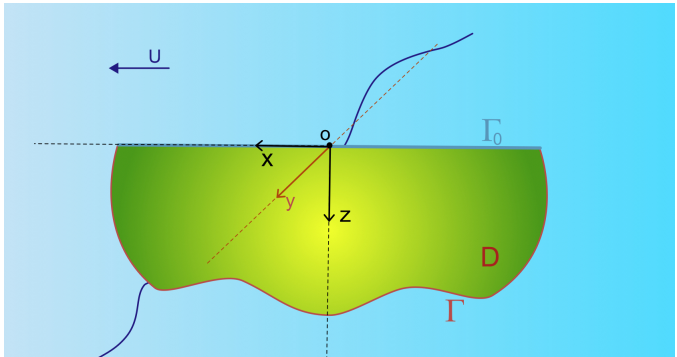
$$\begin{cases} k(U, x, z, x', z') &= \frac{4U^4}{\pi C_F} w(\nu(U)(x - x'), \nu(U)(|z| + |z'|)) \\ w(X, Z) &= \int_1^{\infty} e^{-\lambda^2 Z} \cos(\lambda X) \frac{\lambda^4}{\sqrt{\lambda^2 - 1}} d\lambda \\ \nu(U) &= \frac{g}{U^2} \end{cases}$$

<sup>1</sup>J. Dambrine and M. Pierre, "Regularity of optimal ship forms based on Michell's wave resistance", Appl. Math. Optim. **82**, 23–62 (2020).

# The hull parametrization $f$

The hull parametrization  $f(x, z)$  is computed as a solution to the following problem :

$$\left\{ \text{Find a function } f_D \text{ which minimizes } J_{total}(f) \text{ in the set} \right. \\ \left. \left\{ f : \overline{D} \rightarrow \mathbb{R}, f = 0 \text{ on } \Gamma, f \geq 0 \text{ in } D \text{ and } 2 \int_D f(x, z) dx dz = V \right\} \right\}.$$



## The problem

## Expectation of the water resistance

## Newton's method

## Numerical results

Assume that the variables  $g, \rho$  and  $C_F$  are known physical constants, and  $V$  is the fixed volume of the hull. Given a speed distribution  $U(., \omega)$ , where  $\omega$  is an event in a probability space  $(\mathcal{O}, \mathcal{F}, \mathbb{R})$ , the expectation of the water resistance writes :

$$\begin{aligned}\mathbb{E}(J_{total}) &= \mathbb{E}(J_{viscous} + J_{wave}) \\ &= \int_D |\nabla f(x, z)|^2 dx dz + \int_{\mathcal{O}} \left( \int_{D \times D} k(U(., \omega), x, z, x', z') f(x, z) f(x', z') dx dz dx' dz' \right) d\mathbb{P}(\omega). \\ &= \int_D |\nabla f(x, z)|^2 dx dz + \int_{D \times D} \left( \int_{\mathcal{O}} k(U(., \omega), x, z, x', z') d\mathbb{P}(\omega) \right) f(x, z) f(x', z') dx dz dx' dz'.\end{aligned}$$

i.e the hull parametrisation that minimizes the expectation is solution to :

$$\left\{ \begin{array}{ll} -\Delta f_D(x, z) + \int_D \left( \int_{\mathcal{O}} k(U(., \omega), x, z, x', z') d\mathbb{P} \right) f_D(x', z') dx' dz' &= C \quad (x, z) \in D \\ f_D(x, z) &= 0 \quad \text{on } \Gamma \\ 2 \int_D \hat{f}_D(x, z) dx dz &= V. \end{array} \right.$$

# Numerical results : optimal hulls

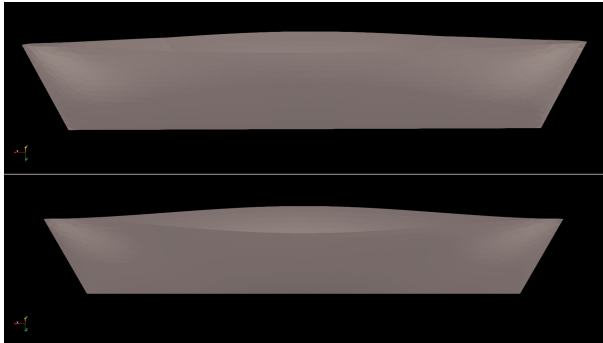
The problem

Expectation of the  
water resistance

Newton's method

Numerical results

- We consider a uniform distribution of velocities  $\nu(\hat{U}(\omega)) \in [0.2\sqrt{gL}, \sqrt{gL}]$ , where  $L = 2$  is the length of the ship. i.e  $Fr = \frac{U}{\sqrt{gL}} \in [0.2, 1]$ .  $D$  is taken as the fixed rectangle of length  $L$  and draft  $T = 0.4$ .



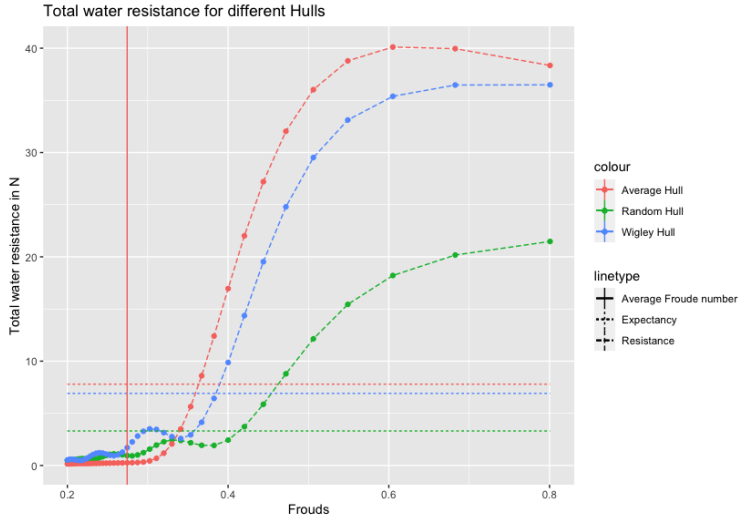
# Numerical results : Total water resistance

The problem

Expectation of the  
water resistance

Newton's method

Numerical results





# Newton's method for the shape optimization problem

The problem

Expectation of the  
water resistance

Newton's method

Numerical results

We now fix an area  $a$  of  $D$ , the optimal design problem reads :

$$\left\{ \begin{array}{l} \text{Find the domain } D^* \text{ which minimizes } \mathbb{E}(J_{total}(f_D)) \\ \text{among all regular open subsets } D \text{ of the lower half-plane} \\ \text{such that } |D| = a \end{array} \right.$$

- A. Novruzi and J. R. Roche, "Newton's method in shape optimisation: a three-dimensional case", *BIT* **40**, 102–120 (2000)
- H. Harbrecht, "A Newton method for Bernoulli's free boundary problem in three dimensions", *Computing* **82**, 11–30 (2008)
- J.-L. Vie, "Second-order derivatives for shape optimization with a level-set method", 2016PESC1072, PhD thesis (2016)

The problem

Expectation of the  
water resistance

Newton's method

Numerical results

Let  $f \in C^2(\mathbb{R}^2; \mathbb{R})$ , and  $E(D) = \int_D f(x) dx$ . We consider the minimization problem

$$\inf_{D \in \mathcal{O}_2} E(D) = D^* = \{x \in \mathbb{R}^N | f(x) < 0\}$$

We search for a descent direction by solving the Newton problem : find  $\theta, \xi \in C^{2,\infty}(\mathbb{R}^N; \mathbb{R}^N)$  such that

$$E''(D; \theta, \xi) = -E'(D; \xi),$$

$$E'(\Omega; \xi) = \int_{\Gamma} (\xi \cdot n) f,$$

$$E''(\Omega; \theta, \xi) = \int_{\Gamma} (\theta \cdot n)(\xi \cdot n)(\mathcal{H}f + \partial_n f) + \int_{\Gamma} Z_{\theta, \xi} f.$$

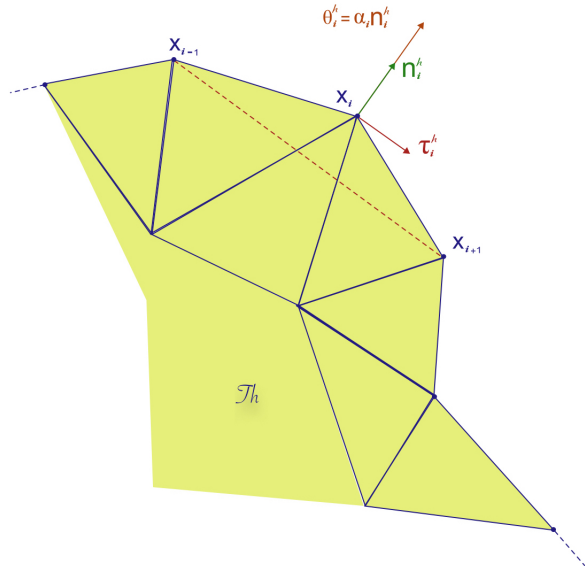
# The proposed discretization

The problem

Expectation of the  
water resistance

Newton's method

Numerical results



# The proposed discretization

The problem

Expectation of the  
water resistance

Newton's method

Numerical results

- In its discrete form, Newton's equation writes as the problem of finding  $\theta^h = (\theta_1^h, \theta_2^h) \in \mathcal{V}_h \times \mathcal{V}_h$  such that  $\forall \xi^h = (\xi_1^h, \xi_2^h) \in \mathcal{V}_h \times \mathcal{V}_h$

$$\int_{\Gamma} (\theta^h \cdot n^h) (\xi^h \cdot n^h) (\mathcal{H}^h f^h + \nabla f^h \cdot n^h) + \int_{\Gamma} Z_{\theta, \xi}^h f^h = - \int_{\Gamma} (\xi^h \cdot n^h) f^h.$$

- If we consider a normal basis, such that  $\theta^h = \sum_{i=0}^{nbe} \alpha_i n_i^h$ , and  $\xi^h = n^h$ , then at every vertex  $x_i$ , a descent direction can be obtained by solving Newton's equation, which is now reduced to

$$\alpha_i \int_{\Gamma} \mathcal{H}_i^h f_i^h + \nabla f_i^h \cdot n_i^h = - \int_{\Gamma} f_i^h \quad \forall i \in \{0, 1, \dots, nbe\}$$

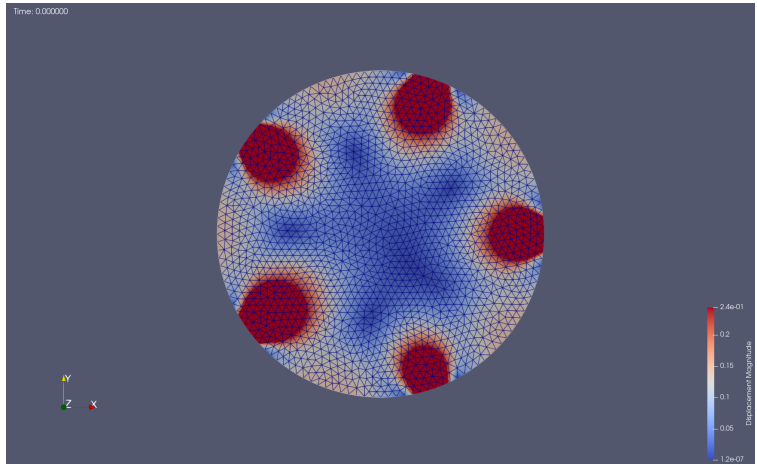
- We thus compute a descent direction on the whole boundary by solving the system  $AT = B$ , where :
  - $A$  is the **diagonal** matrix of size  $nbe \times nbe$ , with entries  $A_{ii} = E''(n_i^h, n_i^h)|_{x_i}$ .
  - $B$  is the vector of size  $nbe$  with entries  $B_i = -E'(n_i^h)|_{x_i}$

The problem

Expectation of the  
water resistance

Newton's method

Numerical results



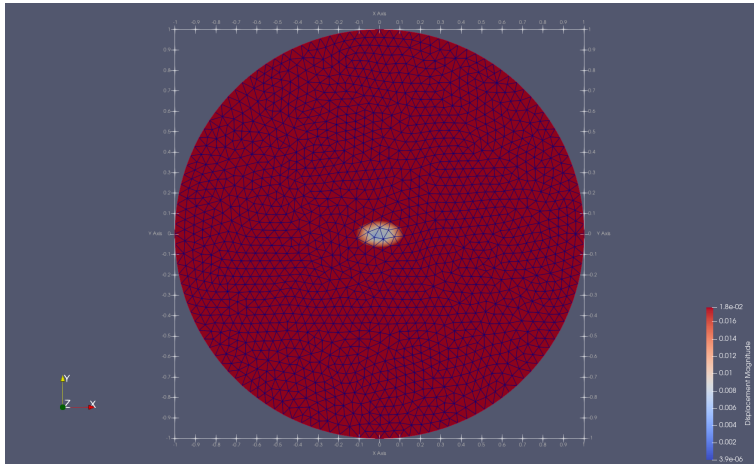
$$f(x, y) = (x^2 + y^2)^5 - 2 \times 0.95^5 (x^5 - 10x^3)y^2 + 5xy^4 + 0.95^{10} - 0.953^{10}$$

The problem

Expectation of the  
water resistance

Newton's method

Numerical results



$$f(x, y) = ((x - 0.5)^2 + y^2)((x + 0.5)^2 + y^2) - 0.51^4$$

The problem

- Consider the set of admissible domains  $\mathcal{A} = \{\Omega : \Omega \subset \mathcal{D}, \Omega \text{ open}, |\Omega| = 1\}$ ,
- Consider the Dirichlet energy :

Expectation of the  
water resistance

$$\begin{aligned} E(\Omega) &= \min \left\{ \frac{1}{2} \int_{\Omega} |\nabla u|^2 dx - \int_{\Omega} u dx : u \in H_0^1(\Omega) \right\} \\ &= -\frac{1}{2} \int_{\Omega} w_{\Omega} dx, \end{aligned}$$

Newton's method

Numerical results

where  $w_{\Omega}$  is the weak solution of the equation  $-\Delta w_{\Omega} = 1 \quad w_{\Omega} \in H_0^1(\Omega)$

- Consider the shape optimization problem  $\min \{E(\Omega) : \Omega \in \mathcal{A}\}$
- to handle the volume constraint we define the Lagrangian :  
 $L(\Omega) = E(\Omega) + \lambda C(\Omega) = E(\Omega) + \lambda(V(\Omega) - V_0)$
- The shape derivative  $w'_{\Omega}(\theta)$  of  $w_{\Omega}$  is solution to the following problem :

$$\begin{cases} -\Delta w'_{\Omega} &= 0 & \text{in } \Omega \\ w'_{\Omega} &= -(\theta \cdot n) \partial_n w_{\Omega} & \text{on } \Gamma \end{cases}$$

The problem

Expectation of the  
water resistance

Newton's method

Numerical results

- The Newton method defines the next iterate  $(\Omega_{k+1}, \lambda_{k+1})$  by :

$$\Omega_{k+1} = (Id + \theta_k)(\Omega_k) \quad \text{and} \quad \lambda_{k+1} = \lambda_k + \mu_k$$

where  $(\theta_k, \mu_k)$  are a solution to

$$\begin{pmatrix} L''_k(\Omega; \theta, \xi) & C'^T(\Omega; \theta) \\ C'(\Omega; \theta) & 0 \end{pmatrix} \begin{pmatrix} \theta_k \\ \mu_k \end{pmatrix} = - \begin{pmatrix} L'_k(\Omega; \xi) \\ c_k \end{pmatrix}$$

- Once again taking  $\theta^h, \xi^h$  normal to the boundary gives the following discretization :

- i  $L''_k(\Omega, n^h, n^h)$  is the diagonal matrix of size  $nbe \times nbe$  with diagonal entries  $L''(\Omega, n_i^h, n_i^h)|_{x^i}$
- ii  $C'(\Omega, n^h)$  is the  $nbe$  vector with entries  $C'(\Omega, n_i^h)|_{x^i}$
- iii  $L'(\Omega, n^h)$  is the  $nbe$  vector with entries  $L'(\Omega, n_i^h)|_{x^i}$
- iv  $c_k$  is the constraint.

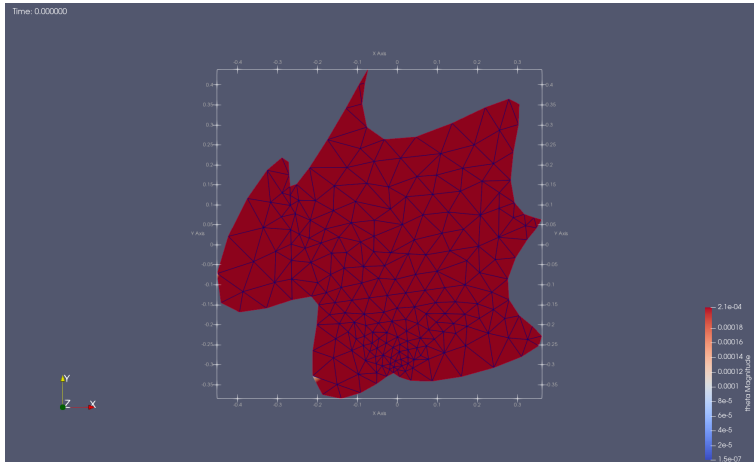


The problem

Expectation of the  
water resistance

Newton's method

Numerical results



minimisation of the Dirichlet energy with volume constraint.

# Back to the water resistance

The problem

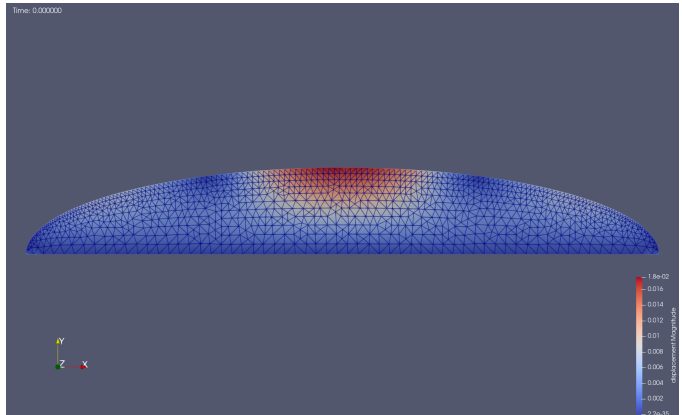
Expectation of the  
water resistance

Newton's method

Numerical results

## Theorem (J.Dambrine, Mo.Pierre)

*Michell's normalized wave resistance kernel  $k$  belongs to  $L^q(D \times D)$  for all  $1 \leq q < \frac{5}{4}$ . Moreover, if  $D$  contains an open disc centered on the  $x$ -axis then  $k$  does not belong to  $L^{5/4}(D \times D)$ .*



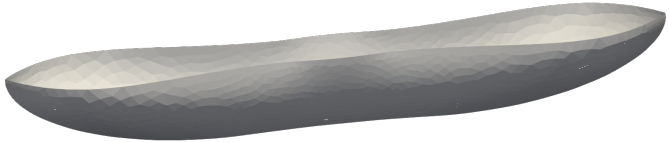
The problem

Expectation of the  
water resistance

Newton's method

Numerical results

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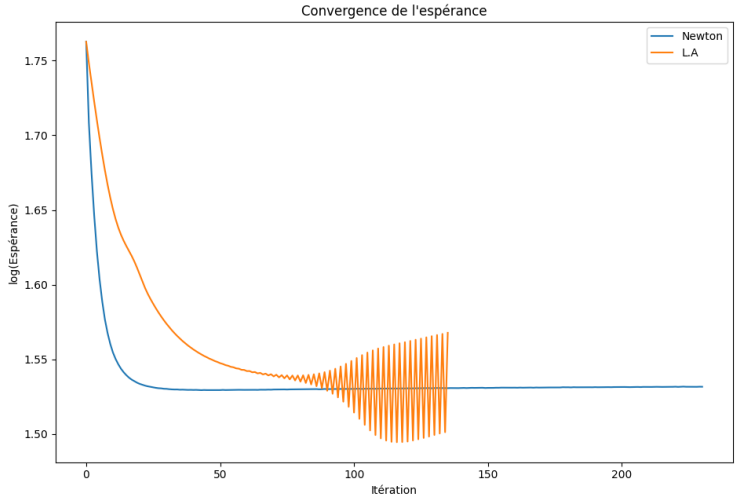
# Numerical results : Convergence of the water resistance

The problem

Expectation of the  
water resistance

Newton's method

**Numerical results**



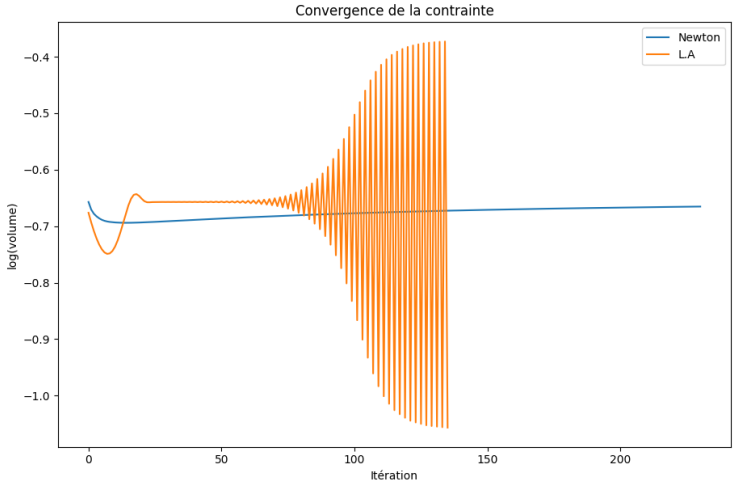
# Numerical results : Convergence of the constraint

The problem

Expectation of the  
water resistance

Newton's method

**Numerical results**



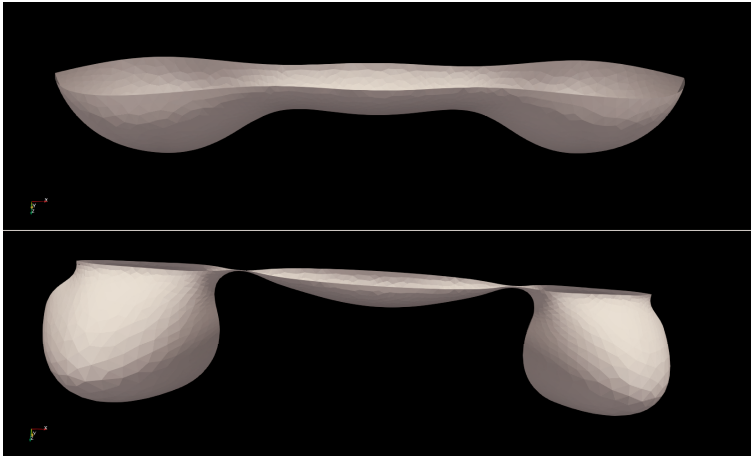
# Numerical results : Optimal hulls

The problem

Expectation of the  
water resistance

Newton's method

**Numerical results**



top : optimal hull for the expectancy, bottom : optimal hull for the average speed.

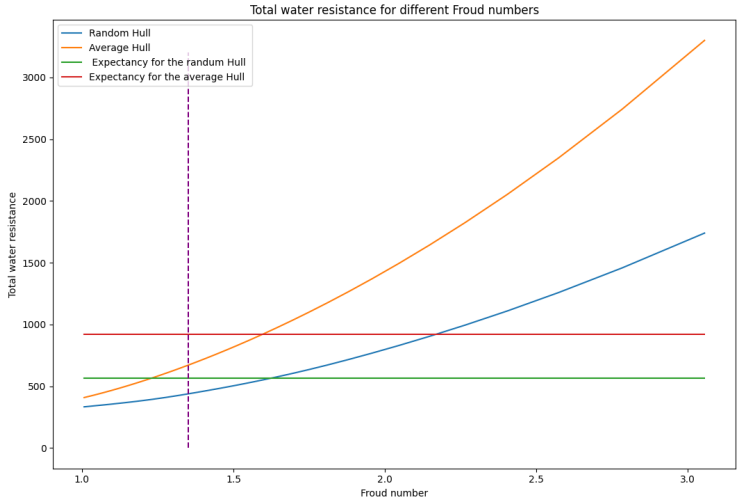
# Numerical results : Comparison for different Froud numbers

The problem

Expectation of the  
water resistance

Newton's method

Numerical results



Thank you

