



Learning the dynamics of systems with memory : Generalized Langevin equations

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Dimensionality reduction

Tackle high dimensional configurational space with a low-dimensional projection.

- Choice of collective variables?
- Effective dynamics?



From all atoms simulations



To reduced dynamical models

- Speed up in computational times (multiscale dynamics)
- Better interpretability (less dimensions)

The MAESTRO project aims at a conceptual and operational breakthrough in computational materials science, particularly in the study of materials for energy applications and storage.



Collaboration between mathematicians, physicists and chemists @ ISCD, including Ludovic Goudenège, Pierre Monmarché, Fabio Pietrucci and Benjamin Rotenberg. Prediciting thermodynamics (free energy) and kinetics (rates, timescales) properties of energy storage materials from numerical simulations.

1) The Generalized Langevin Equation (GLE): here comes the memory

2) Infering the dynamics?

3) Expectation maximization algorithm

4) Examples

The Generalized Langevin Equation (GLE): here comes the memory





Dynamical (Hamilton) Equations:

$$\begin{cases} \dot{x}_{\text{red}} = v_{\text{red}} \\ \dot{v}_{\text{red}} = F_{r \to r}(x_{\text{red}}) + F_{g \to r}(x_{\text{red}}, x_{\text{green}}) \end{cases} \begin{cases} \dot{x}_{\text{green}} = v_{\text{green}} \\ \dot{v}_{\text{green}} = F_{g \to g}(x_{\text{green}}) + F_{r \to g}(x_{\text{red}}, x_{\text{green}}) \end{cases}$$

Let's focus on red particles as collective variables.

Can we remove green particles variables? Yes, using Mori Zwanzig projection formalism.



Generalized Langevin equation¹:

 $\dot{x}_{red} = v_{red}$

$$\dot{v}_{\mathsf{red}} = \widetilde{F}(x_{\mathsf{red}}) - \int_0^t K(\tau) v_{\mathsf{red}}(t-\tau) \mathrm{d}\tau + \xi(t)$$



The solvant is modeled via

- a non-Markovian memory (friction) kernel K(τ) (dissipation of energy)
- a diffusion term ξ(t) (random noise coming from collision). This is a function of initial conditions

¹Zwanzig, 2001

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Position-dependent form of the Generalized Langevin Equations for Hamiltonian $\mathsf{systems}^2$:

$$\dot{v}(t) = \widetilde{F}(x(t)) - \int_0^t K(\tau, x(t-\tau))v(t-\tau)\mathrm{d} au + \xi(t)$$

¹Zwanzig, 2001 ²Vroylandt and Monmarché 2022, arXiv:2201.02457

What are we modeling with this memory? Dynamical timescales of the rest of the system



We can have "short" or "long" memory. We focus on "short" memory. $K(\tau) = \delta(\tau) \Leftrightarrow$ Langevin description

Infering the dynamics?

The inference question

You are given:



How do you infer the parameters in this equation from simulation data?

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How do you infer the parameters in this equation from simulation data?

How do you generate new trajectories from our inferred dynamics



State of the art for memory fitting

Most of current approaches derive from GLE

$$\dot{\mathbf{v}} = \widetilde{F}(\mathbf{x}) - \int_0^t K(\tau) \mathbf{v}(t- au) \mathrm{d} au + \xi(t)$$

a Volterra integral equation for the memory kernel

GLE

$$\dot{v} = \widetilde{F}(x) - \int_{0}^{t} K(\tau)v(t-\tau)d\tau + \xi(t)$$
a Volterra integral equation for the memory
kernel

$$\left\langle v(0)\left(\dot{v}(t) - \widetilde{F}(x)\right)\right\rangle = -\int_{0}^{t} K(\tau)\left\langle v(0)v(t-\tau)\right\rangle d\tau$$

$$\tau$$
Volterra 2nd kind
True Kernel

$$\left\langle v(0)\left(\dot{v}(t) - \widetilde{F}(x)\right)\right\rangle = -\int_{0}^{t} K(\tau)\left\langle v(0)v(t-\tau)\right\rangle d\tau$$

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- Laplace transform¹
- Integral discretization²
- ...

^aLei, Baker & Li, PNAS 2016, 113 ^bDaldrop, Kappler, Brünig, & Netz PNAS, 2018, 115

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angle \mathrm{d} au$$

 $K(\tau)$

That can be solved using

- Laplace transform¹
- Integral discretization²

• ...

But

- This does not give a generative model
- Convergence issue for some cases



^aLei, Baker & Li, PNAS 2016, 113 ^bDaldrop, Kappler, Brünig, & Netz PNAS, 2018, 115

 τ

And the noise?

The noise is obtained from the inversion of

$$\dot{x} = v$$
$$\dot{v} = F(x) - \int_0^t K(\tau) v(t - \tau) d\tau + \xi(t)$$

It exhibits non Gaussian tails and correlation in time



Auxiliary/hidden variables approach

Generalized Langevin equation:

$$\begin{cases} \dot{x} = v \\ \dot{v} = F(x) - \int_0^t K(\tau) v(t-\tau) d\tau + \xi(t) \end{cases} \begin{cases} \dot{x} = v \\ \dot{v} = F(x) - A_{vh} h - A_{vv} v + \sigma_v W_v(t) \\ \dot{h} = -A_{hh} h - A_{hv} v + \sigma_h W_h(t) \end{cases}$$

Aiming to capture most important missed degrees of freedom.

Auxiliary/hidden variables approach

Generalized Langevin equation:

$$\begin{cases} \dot{x} = v \\ \dot{v} = F(x) - \int_0^t K(\tau) v(t-\tau) d\tau + \xi(t) \end{cases}$$

Hidden variable

$$\begin{cases}
\dot{x} = v \\
\dot{v} = F(x) - A_{vh} \dot{h} - A_{vv} v + \sigma_v W_v(t) \\
\dot{h} = -A_{hh} \dot{h} - A_{hv} v + \sigma_h W_h(t)
\end{cases}$$

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Gaussian white noise

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Aiming to capture most important missed degrees of freedom.

Equivalent equations? Formal solution to last equation:

$$\boldsymbol{h}(t) = \int_0^t e^{-\boldsymbol{A}_{hh}\tau} \boldsymbol{A}_{hv} \boldsymbol{v}(t-\tau) \mathrm{d}\tau + \int_0^t e^{-\boldsymbol{A}_{hh}\tau} \boldsymbol{\sigma}_h \mathrm{d} \boldsymbol{W}_h(\tau)$$

Memory kernel is the sum of a Dirac and a Prony series:

$$K(\tau) = A_{vv}\delta(\tau) + \sum_{i}^{\dim h} c_i e^{-\lambda_i \tau}$$

Generalized Langevin equation:

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Aiming to capture most important missed degrees of freedom.



- GLE with memory kernel \Rightarrow Non-markovian but explicit
- Langevin equations with auxiliary variables \Rightarrow Markovian but hidden parts

The inference question (bis)

You are given:



Model: Generalized Langevin equation with hidden variables

$$\begin{cases} \dot{\mathbf{x}} = \mathbf{v} \\ \dot{\mathbf{v}} = \mathbf{F}(\mathbf{x}) - \mathbf{A}_{vh}\mathbf{h} - \mathbf{A}_{vv}\mathbf{v} + \sigma_{v}\mathbf{W}_{v}(t) \\ \dot{\mathbf{h}} = -\mathbf{A}_{hh}\mathbf{h} - \mathbf{A}_{hv}\mathbf{v} + \sigma_{h}\mathbf{W}_{h}(t) \end{cases}$$

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How do you generate new trajectories from our inferred dynamics

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How do you infer the parameters in this equation from simulation data? Expectation maximization algorithm.

- Iterative algorithm to deal with hidden variables
- E-step: expectation (recovering hidden variables history)
- M-step: maximization of likelihood (recovering parameters)

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How do you generate new trajectories from our inferred dynamics

Once the parameters are fitted, this is a generative model.

Expectation maximization algorithm

$$\begin{cases} \dot{\mathbf{x}} = \mathbf{v} \\ \dot{\mathbf{v}} = \mathbf{F}(\mathbf{x}) - \mathbf{A}_{vh}\mathbf{h} - \mathbf{A}_{vv}\mathbf{v} + \sigma_{v}\mathbf{W}_{v}(t) \\ \dot{\mathbf{h}} = -\mathbf{A}_{hh}\mathbf{h} - \mathbf{A}_{hv}\mathbf{v} + \sigma_{h}\mathbf{W}_{h}(t) \end{cases}$$

Discretizing the dynamics

We use a numerical scheme that discretizes the Langevin equation and a parametrization $\Theta = \{F, A, \sigma\}$ of the (force, friction, diffusion) coefficients.

$$(\boldsymbol{X}_{t+\Delta t}, \boldsymbol{h}_{t+\Delta t}) = (\boldsymbol{X}_t, \boldsymbol{h}_t) + G_{\Theta}(\boldsymbol{X}_t, \boldsymbol{h}_t) + \boldsymbol{S}_{\Theta} \boldsymbol{W}(t)$$

We separate between visible X_t and hidden h_t variables

Transition probability

Hence we have a gaussian transition probability

$$\mathbb{P}\left(\boldsymbol{X}_{t+\Delta t}, \boldsymbol{h}_{t+\Delta t} | \boldsymbol{X}_{t}, \boldsymbol{h}_{t}\right)_{\boldsymbol{\Theta}} = \mathcal{N}\left(\left(\boldsymbol{X}_{t}, \boldsymbol{h}_{t}\right) + G_{\boldsymbol{\Theta}}(\boldsymbol{X}_{t}, \boldsymbol{h}_{t}), \boldsymbol{S}_{\boldsymbol{\Theta}}\boldsymbol{S}_{\boldsymbol{\Theta}}^{\mathrm{T}}\right)$$

E-step: Reconstructing hidden variables

If you know the parameters: Markovianity of the system \Rightarrow We use a predictor-corrector-smoother (or Kalman filer and Rauch smoother) Using evolution equation

$$\begin{pmatrix} \boldsymbol{X}_{t+\Delta t} \\ \boldsymbol{h}_{t+\Delta t} \end{pmatrix} = \begin{pmatrix} \boldsymbol{X}_t \\ \boldsymbol{h}_t \end{pmatrix} + \boldsymbol{G}_{\boldsymbol{\Theta}}(\boldsymbol{X}_t, \boldsymbol{h}_t) + \boldsymbol{S}_{\boldsymbol{\Theta}} \boldsymbol{W}(t)$$

You get probabilistic reconstruction of the hidden variables $\mathbb{P}\left(\{\boldsymbol{h}_t\}_0^T | \{\boldsymbol{X}_t\}_0^T, \boldsymbol{\Theta}\right)$





$$\mathbb{P}\left(\{\boldsymbol{X}_{t},\boldsymbol{h}_{t})\}_{0}^{T}|\boldsymbol{\Theta}\right) = \mathbb{P}\left((\boldsymbol{X}_{0},\boldsymbol{h}_{0})\right)\prod_{t=0}^{T}\mathbb{P}\left((\boldsymbol{X}_{t+\Delta t},\boldsymbol{h}_{t+\Delta t}|(\boldsymbol{X}_{t},\boldsymbol{h}_{t}))_{\boldsymbol{\Theta}}\right)$$

When dealing with hidden variables, we average over possible hidden variables values

$$\Theta^* = \arg \max_{\Theta} \int d\{\boldsymbol{h}_t\}_0^T \mathbb{P}\left(\{\boldsymbol{h}_t\}_0^T | \{\boldsymbol{X}_t\}_0^T, \Theta\right) \log \mathcal{L}(\{\boldsymbol{X}_t, \boldsymbol{h}_t\}_0^T; \Theta)$$

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Algorithm - Iterative strategy

• Initialize parameters at some value Θ_0

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- M-step: Maximize the expectation of the log-likelihood
- Iterate until convergence



Examples

- Quadratic potential
- 5 hidden dimensions
- Random value for the friction matrix A with an hierarchical structure.
- 20 trajectories of 25.10³ timesteps.



Validation of the algorithm

- Quadratic potential
- 5 hidden dimensions
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Non-equilibrium 2D setup¹

- Quadratic potential with two different temperatures in *x* and *y* directions.
- 2 hidden dimensions
- Random value for the friction matrix A.
- 20 trajectories of 30.10³ timesteps.



¹Mancois, Marcos, Viot, and Wilkowski, 2018

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¹Mancois, Marcos, Viot, and Wilkowski, 2018

A molecular dynamics example. We set up 512 LJ particles in a 3D box with periodic boundary conditions. 2 of the particles form a dimer.



Distance r between the two dimer particles as collective variable.

Lennard-Jones Dimer



Generate 50 trajectories of 2×100000 timesteps from the fitted model.

We can change the timestep if wanted.

First passage time



- This method fits the full statistical model for the dynamics (force + friction + noise).
- Generative model \Rightarrow Possibility to use the inferred dynamics in future computations.
- Works for multidimensionnal and non equilibrium systems.
- Can be quite long to run...

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Thank you for your attention and questions?

Vroylandt, Hadrien, Ludovic Goudenège, Pierre Monmarché, Fabio Pietrucci, and Benjamin Rotenberg. *Likelihood-based non-Markovian models from molecular dynamics.* PNAS, 2022, 119

Hadrien Vroylandt and Pierre Monmarché. Position-dependent memory kernel in generalized Langevin equations: theory and numerical estimation. arXiv:2201.02457 (accepted in Journal of Chemical Physics) Let's assume that our green particles are harmonic particles, harmonically coupled to the red ones.

$$\begin{cases} \dot{x}_{\text{red}} = v_{\text{red}} \\ \dot{v}_{\text{red}} = F_{r \to r}(x_{\text{red}}) + \sum_{j} \gamma_j (x_{\text{green}}^j - \frac{\gamma_j}{\omega_j^2} x_{\text{red}}) \end{cases} \begin{cases} \dot{x}_{\text{green}}^j = v_{\text{green}}^j \\ \dot{v}_{\text{green}}^j = -\omega_j^2 x_{\text{green}}^j + \gamma_j x_{\text{red}} \end{cases}$$

The motion of the green particles is easily solved (Second order differential equation + integration by part):

$$x_{\text{green}}^{j}(t) - \frac{\gamma_{j}}{\omega_{j}^{2}} x_{\text{red}}(t) = x_{\text{green}}^{j}(0) \cos \omega_{j} t - \gamma_{j} \int_{0}^{t} v_{\text{red}}(t-s) \frac{\cos \omega_{j}(s)}{\omega_{j}} ds$$

This lead to the equation of motion for the red particules

$$\begin{cases} \dot{x}_{\text{red}} = v_{\text{red}} \\ \dot{v}_{\text{red}} = F_{r \to r}(x_{\text{red}}) - \int_0^t v_{\text{red}}(t-s) \sum_j \frac{\cos \omega_j(s)}{\omega_j} ds + \sum_j x_{\text{green}}^j(0) \cos \omega_j t \end{cases}$$

Projection operator

For a given basis function family $\{b\}$, we define the projector as

$$\mathcal{P}_{\{b\}}F = B \cdot \langle B, B \rangle^{-1} \cdot \langle B, F \rangle$$

with scalar product $\langle f,g\rangle=\int \mathrm{d} \pmb{X} \rho_{eq}(\pmb{X}) f^{\mathrm{T}}(\pmb{X}) g(\pmb{X})$

The evolution equation of our observables are given by

Evolution equation

$$\frac{\partial \mathcal{O}(\boldsymbol{X},t)}{\partial t} = \mathcal{LO}(\boldsymbol{X},t) = e^{t\mathcal{L}}\mathcal{LO}(\boldsymbol{X},0) = e^{t\mathcal{L}}\mathcal{P}_{\{b\}}\mathcal{LO}(\boldsymbol{X},0) + e^{t\mathcal{L}}(1-\mathcal{P}_{\{b\}})\mathcal{LO}(\boldsymbol{X},0)$$

Duhamel-Dyson identity: $e^{t\mathcal{L}} = e^{t(1-\mathcal{P}_{\{b\}})\mathcal{L}} + \int \mathrm{d}s e^{(t-s)\mathcal{L}} \mathcal{P}_{\{b\}}\mathcal{L}e^{s(1-\mathcal{P}_{\{b\}})\mathcal{L}}$

$$\begin{split} \frac{\partial \mathcal{O}(\boldsymbol{X},t)}{\partial t} = & e^{t\mathcal{L}} \mathcal{P}_{\{b\}} \mathcal{L} \mathcal{O}(\boldsymbol{X},0) + \int \mathrm{d} s e^{(t-s)\mathcal{L}} \mathcal{P}_{\{b\}} \mathcal{L} e^{s(1-\mathcal{P}_{\{b\}})\mathcal{L}} (1-\mathcal{P}_{\{b\}}) \mathcal{L} \mathcal{O}(\boldsymbol{X},0) \\ &+ e^{t(1-\mathcal{P}_{\{b\}})\mathcal{L}} (1-\mathcal{P}_{\{b\}}) \mathcal{L} \mathcal{O}(\boldsymbol{X},0) \\ &\frac{\partial \mathcal{O}(\boldsymbol{X},t)}{\partial t} = e^{t\mathcal{L}} \mathcal{P}_{\{b\}} \mathcal{L} \mathcal{O}(\boldsymbol{X},0) + \int \mathrm{d} s e^{(t-s)\mathcal{L}} \mathcal{P}_{\{b\}} \mathcal{L} F(\boldsymbol{X},s) + F(\boldsymbol{X},t) \\ &\text{with } F(\boldsymbol{X},t) = e^{t(1-\mathcal{P}_{\{b\}})\mathcal{L}} (1-\mathcal{P}_{\{b\}}) \mathcal{L} \mathcal{O}(\boldsymbol{X},0) \end{split}$$