Analysis of a domain decomposition method for the convected Helmholtz equation

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In this work, we are interested in solving an Helmholtz like equation:

 $-\operatorname{div}(A \nabla \mathbf{U}) + i\mathbf{a} \cdot \nabla \mathbf{U} + \mu \mathbf{U} = \mathbf{f}$

where: * A is a 2x2 symmetric positive definite matrix,

- * a is a vecteur of \mathbb{R}^2 ,
- * μ is a real constant.

In this work, we are interested in solving an Helmholtz like equation:

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This type of equation occurs in several contexts:

* The convected Helmholtz equation $(A = c_0 I - vv^t, a = -2\omega v, \mu = -\omega^2)$

$$-\operatorname{div}((c_0 - \mathbf{v}\mathbf{v}^t) \nabla \mathbf{u}) - 2i\omega\mathbf{v} \cdot \nabla \mathbf{u} + \omega^2 \mathbf{u} = \mathbf{f}$$

H. Barucq et al, HDG and HDG+ methods for harmonic wave problems with convections, 2021



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* The Gröss-Pitaevskii equation (computation of the ground states)

I. Danaila et al, Computation of ground states of the Gröss-Pitaevskii functional via Riemannian optimization, 2017

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* The Gröss-Pitaevskii equation (computation of the ground states)

- * The wave-ray equation (A = I, a = v, $\mu = 0$)
 - $-\Delta \mathbf{U} + i\mathbf{a} \cdot \nabla \mathbf{U} = \mathbf{f}$

P. Verburg et al, Multi-level wave-ray method for 2d Helmholtz equation, 2010

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Goal: Propose an efficient iterative algorithm of resolution

In short, it is as difficult as solving the Helmholtz equation !!



O.G. Ernst et al, Why is it difficult to solve the Helmholtz equation? 2012



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1. Motivation

- 2. Link with Helmholtz equation
- 3. Convergence analysis on a toy problem
- 4. An alternated iterative algorithm
- 5. Conclusion

Let us consider u solution to

 $-\operatorname{div}(A \nabla \mathbf{U}) + i\mathbf{a} \cdot \nabla \mathbf{U} + \mu \mathbf{U} = \mathbf{f}$

Then, setting $\mathbf{u} = e^{\imath k \cdot x} \mathbf{v}'$ with $k = \frac{1}{2} \mathbf{A}^{-1} \mathbf{a}$ one get that

$$\operatorname{div}(A \nabla \mathbf{v}') + \left(\mu - \frac{\|\mathbf{a}\|_{A^{-1}}^2}{4}\right) \mathbf{v}' = \mathbf{f}'$$



Remark: Even if $\mu \ge 0$, we see that the problem is not coercive if $||\mathbf{a}||$ is large.

Let us consider u solution to

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Now, taking the change of variables $(x, y) \leftarrow T(x, y)$ where T is a matrix, we get

$$-\operatorname{div}(TAT^{t}\nabla\mathbf{v}) + \left(\mu - \frac{\|\mathbf{a}\|_{A^{-1}}^{2}}{4}\right)\mathbf{v} = \widetilde{\mathbf{f}}'$$

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A simple idea then to obtain the Helmholtz equation is to take $T = G^{-1}$ where $A = GG^{t}$.

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A simple idea then to obtain the Helmholtz equation is to take $T = G^{-1}$ where $A = GG^{t}$.

Remark: The choice of the transformation is not unique!



F.Q. Hu et al, On the use of Prandtl-Glauert-Lorentz transformation for acoustic scattering by rigid bodies with a uniform flow, 2019

Y. Gao et al, Wave scattering in layered orthotropic media I: a stable PML and a highaccuracy boundary integral equation method, 2021

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Cartesian PML formulation:

 $-\operatorname{div}(A \nabla \mathbf{U}) + i\mathbf{a} \cdot \nabla \mathbf{U} + \mu \mathbf{U} = \mathbf{f}$

Convected Helmholtz equation

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Convected Helmholtz equation

Coordinates transformation

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 $-\Delta \mathbf{V} + \widetilde{\omega} \mathbf{U} = \mathbf{f}$

Helmholtz equation

Cartesian PML formulation:

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Convected Helmholtz equation

Coordinates transformation

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$$-\Delta \mathbf{v} + \widetilde{\omega} \mathbf{u} = \mathbf{f}$$

Helmholtz equation

Complex stretching

$$-\operatorname{div}(D_{PML}\nabla \mathbf{v}_{PML}) + \rho_{PML}\widetilde{\omega} \mathbf{u}_{PML} = \mathbf{f}$$

PML Helmholtz equation

Cartesian PML formulation:

 $-\operatorname{div}(A \nabla \mathbf{U}) + i\mathbf{a} \cdot \nabla \mathbf{U} + \mu \mathbf{U} = \mathbf{f}$ Coordinates transformation $-\Delta \mathbf{V} + \widetilde{\omega} \mathbf{U} = \mathbf{f}$ Helmholtz equation $-\operatorname{div}(A_{PML} \nabla \mathbf{U}_{PML}) + \frac{i}{2} \mathbf{a}_{PML} \cdot \nabla \mathbf{U}_{PML}$ $+ \frac{i}{2} \operatorname{div}(\mathbf{a}_{PML} \mathbf{U}_{PML}) + \mu_{PML} \mathbf{U}_{PML} = \mathbf{f}$ PML Convected Helmholtz equation $\operatorname{Coordinates transformation} -\Delta \mathbf{V} + \widetilde{\omega} \mathbf{U} = \mathbf{f}$ Helmholtz equation $-\operatorname{div}(D_{PML} \nabla \mathbf{V}_{PML}) + \rho_{PML} \widetilde{\omega} \mathbf{U}_{PML} = \mathbf{f}$ PML Convected Helmholtz equation $\operatorname{Coordinates transformation} PML Helmholtz equation$



P. Marchner et al, Stable Perfectly Matched Layers with Lorentz transformation for the convected Helmholtz equation, 2019

E. Becache et al, Perfectly matched layers for the convected Helmholtz equation, 2004

Cartesian PML formulation:

 $-\operatorname{div}(A \nabla \mathbf{U}) + i\mathbf{a} \cdot \nabla \mathbf{U} + \mu \mathbf{U} = \mathbf{f}$

Convected Helmholtz equation

Coordinates transformation

$$\Delta \mathbf{v} + \widetilde{\omega} \mathbf{u} = \mathbf{f}$$

Helmholtz equation

Complex stretching

$$\operatorname{div}(A_{PML} \nabla \mathbf{U}_{PML}) + \frac{\iota}{2} \mathbf{a}_{PML} \cdot \nabla \mathbf{U}_{PML}$$

 $+\frac{\iota}{2}\mathrm{div}(\mathbf{a}_{PML}\mathbf{U}_{PML})+\mu_{PML}\mathbf{U}_{PML}=\mathbf{f}$

Coordinates transformation

 $-\operatorname{div}(D_{PML} \nabla \mathbf{v}_{PML}) + \rho_{PML} \widetilde{\omega} \mathbf{u}_{PML} = \mathbf{f}$ $\mathbf{PML} \text{ Helmholtz equation}$

PML Convected Helmholtz equation



Illustration (Convected Hemlholtz): $\mathbf{a} = 2\omega \mathbf{v}, \quad \omega = 20$ $\mathbf{v} = [0.8, \ 0]^t$ $A = Id - \mathbf{v}\mathbf{v}^t$

Cartesian PML formulation:

 $-\operatorname{div}(A \nabla \mathbf{U}) + i\mathbf{a} \cdot \nabla \mathbf{U} + \mu \mathbf{U} = \mathbf{f}$ Convected Helmholtz equation $-\operatorname{div}(A_{PML} \nabla \mathbf{U}_{PML}) + \frac{i}{2} \mathbf{a}_{PML} \cdot \nabla \mathbf{U}_{PML}$ $+ \frac{i}{2} \operatorname{div}(\mathbf{a}_{PML} \mathbf{U}_{PML}) + \mu_{PML} \mathbf{U}_{PML} = \mathbf{f}$ PML Convected Helmholtz equation







Cartesian PML formulation:



ABC (Absorbing Boundary Conditions):

N. Rouxelin et al, Prandtl-Glauert-Lorentz based Absorbing Boundary Conditions for the convected Helmholtz equation, 2021



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Formulation of the problem

 $-\operatorname{div}(A \nabla \mathbf{U}) + i\mathbf{a} \cdot \nabla \mathbf{U} + \mu \mathbf{U} = \mathbf{f} \quad \text{in} \quad \Omega$ $\mathbf{U} = 0 \qquad \qquad \text{on} \quad \partial \Omega$

We will consider four configurations :



Formulation of the problem

-($liv(A \nabla \mathbf{U}) +$	ia ·	VU	+ μ υ	= f	in	Ω
U =	= 0					on	$\partial \Omega$

We will consider four configurations :



In each case, we will consider a Schwarz iterative algorithm of resolution with 2 subdomains.

$$\begin{aligned} \mathscr{L}_{CH} \mathbf{U}^{1,n} &= \mathbf{f}_1 & \text{in } \Omega_1 \\ \mathbf{U}^{1,n} &= 0 & \text{on } \partial\Omega \\ (\partial_n + p_{1,2}) \mathbf{U}^{1,n} &= (\partial_n + p_{1,2}) \mathbf{U}^{2,n-1} & \text{on } \Gamma_{12} \end{aligned}$$

Formulation of the problem

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$\mathscr{L}_{CH} \mathbf{U}^{1,n} = \mathbf{f}_1$	in	Ω ₁	$\mathscr{L}_{CH} \mathbf{U}^{12,n} = \mathbf{f}_2$	in	Ω ₂
$\mathbf{U}^{1,n} = 0$	on	$\partial \Omega$	$\mathbf{U}^{2,n} = 0$	on	λΩ
$(\partial_n + p_{1,2})\mathbf{U}^{1,n} = (\partial_n + p_{1,2})\mathbf{U}^{2,n-1}$	on	Γ ₁₂	$(\partial_n + p_{2,1})\mathbf{U}^{2,n} = (\partial_n + p_{2,1})\mathbf{U}^{1,n-1}$	on	Γ ₂₁

Schwarz algorithm:

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Equivalent Schwarz algorithm for Helmholtz equation:

$$\begin{aligned} & \mathscr{L}_{H} \mathbf{v}^{1,n} = \widetilde{\mathbf{f}}_{1} & \text{in } \widetilde{\Omega}_{1} & \mathscr{L}_{H} \mathbf{v}^{2,n} = \widetilde{\mathbf{f}}_{2} & \text{in } \widetilde{\Omega}_{2} \\ & \mathbf{v}^{1,n} = 0 & \text{on } \partial \widetilde{\Omega} & \mathbf{v}^{2,n} = 0 & \text{on } \partial \widetilde{\Omega} \\ & (\partial_{x} + \widetilde{p}_{1,2}) \mathbf{v}^{1,n} = (\partial_{x} + \widetilde{p}_{1,2}) \mathbf{v}^{2,n-1} & \text{on } \widetilde{\Gamma}_{12} & (\partial_{x} + \widetilde{p}_{2,1}) \mathbf{v}^{2,n} = (\partial_{x} + \widetilde{p}_{2,1}) \mathbf{v}^{1,n-1} & \text{on } \widetilde{\Gamma}_{21} \end{aligned}$$

Schwarz algorithm:

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Remarks: * The convergence analysis can be done only for the Helmholtz equation.

Schwarz algorithm:

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Equivalent Schwarz algorithm for Helmholtz equation:

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Remarks: * The convergence analysis can be done only for the Helmholtz equation.

* To preserve the separable geometry in the Helmholtz case, we need to assume that A is diagonale.

Schwarz algorithm:



Equivalent Schwarz algorithm for Helmholtz equation:

$$\begin{aligned} \mathscr{L}_{H} \mathbf{v}^{1,n} &= \widetilde{\mathbf{f}}_{1} & \text{in } \widetilde{\Omega}_{1} & \mathscr{L}_{H} \mathbf{v}^{2,n} &= \widetilde{\mathbf{f}}_{2} & \text{in } \widetilde{\Omega}_{2} \\ \mathbf{v}^{1,n} &= 0 & \text{on } \partial \widetilde{\Omega} & \mathbf{v}^{2,n} &= 0 & \text{on } \partial \widetilde{\Omega} \\ (\partial_{x} + \widetilde{p}_{1,2}) \mathbf{v}^{1,n} &= (\partial_{x} + \widetilde{p}_{1,2}) \mathbf{v}^{2,n-1} & \text{on } \widetilde{\Gamma}_{12} & (\partial_{x} + \widetilde{p}_{2,1}) \mathbf{v}^{2,n} &= (\partial_{x} + \widetilde{p}_{2,1}) \mathbf{v}^{1,n-1} & \text{on } \widetilde{\Gamma}_{21} \end{aligned}$$

Remarks: * Optimized TC can be derived from optimized TC for the Helmholtz equation

M.J. Gander et al, Optimized schwarz methods with overlap for the helmholtz equation 2016

Convergence analysis (D-D case)



Remark: For simplicity, we will assume that A = Id s.t. $\widetilde{\Omega} = [0,1]^2$, $\widetilde{\Gamma}_{1,2} = \{\beta\} \times [0,1]$ and $\widetilde{\Gamma}_{2,1} = \{\beta\} \times [0,1]$ 14

Convergence analysis (D-D case)



Using separation of variables methods, we get that

$$\mathbf{v}^{i,n} = \sum_{\xi \in \mathbb{N}} N_{\xi} \sin(\xi \pi y) \left(\mathbf{A}^{i,n}(\xi) e^{i\mathcal{S}(\xi)x} + \mathbf{B}^{i,n}(\xi) e^{-i\mathcal{S}(\xi)x} \right) \quad \text{where } \mathcal{S}(\xi) = \sqrt{\widetilde{\mu} - (\xi \pi)^2}.$$

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The boundary conditions implies: $\mathbf{B}^{1,n}(\xi) = -\mathbf{A}^{1,n}(\xi)$ and $\mathbf{B}^{2,n}(\xi) = -e^{2i\mathcal{S}(\xi)}\mathbf{A}^{2,n}(\xi)$

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$$\begin{array}{ll} \text{The TC implies:} \quad \mathbf{A}^{1,n}(\xi) = \frac{e^{i\mathcal{S}(\xi)\beta}(i\mathcal{S}(\xi) + \widetilde{p}_{1,2}) + e^{2i\mathcal{S}(\xi)}e^{-i\mathcal{S}(\xi)\beta}(i\mathcal{S}(\xi) - \widetilde{p}_{1,2})}{e^{i\mathcal{S}(\xi)\beta}(i\mathcal{S}(\xi) + \widetilde{p}_{1,2}) + e^{-i\mathcal{S}(\xi)\beta}(i\mathcal{S}(\xi) - \widetilde{p}_{1,2})} \mathbf{A}^{2,n-1}(\xi) \\ \\ \mathbf{A}^{2,n}(\xi) = \frac{e^{i\mathcal{S}(\xi)\alpha}(i\mathcal{S}(\xi) + \widetilde{p}_{2,1}) + e^{-i\mathcal{S}(\xi)\alpha}(i\mathcal{S}(\xi) - \widetilde{p}_{2,1})}{e^{i\mathcal{S}(\xi)\alpha}(i\mathcal{S}(\xi) + \widetilde{p}_{2,1}) + e^{2i\mathcal{S}(\xi)e^{-i\mathcal{S}(\xi)\alpha}(i\mathcal{S}(\xi) - \widetilde{p}_{2,1})} \mathbf{A}^{1,n-1}(\xi) \end{array}$$

Convergence analysis (D-D case)

$$\begin{aligned} & \mathscr{L}_{H} \mathbf{v}^{1,n} = 0 & \text{in } \widetilde{\Omega}_{1} & \mathscr{L}_{H} \mathbf{v}^{2,n} = 0 & \text{in } \widetilde{\Omega}_{2} \\ & \mathbf{v}^{1,n} = 0 & \text{on } \partial \widetilde{\Omega} & \mathbf{v}^{2,n} = 0 & \text{on } \partial \widetilde{\Omega} \\ & (\partial_{x} + \widetilde{p}_{1,2}) \mathbf{v}^{1,n} = (\partial_{x} + \widetilde{p}_{1,2}) \mathbf{v}^{2,n-1} & \text{on } \widetilde{\Gamma}_{12} & (\partial_{x} + \widetilde{p}_{2,1}) \mathbf{v}^{2,n} = (\partial_{x} + \widetilde{p}_{2,1}) \mathbf{v}^{1,n-1} & \text{on } \widetilde{\Gamma}_{21} \end{aligned}$$

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The boundary conditions implies: $\mathbf{B}^{1,n}(\xi) = -\mathbf{A}^{1,n}(\xi)$ and $\mathbf{B}^{2,n}(\xi) = -e^{2i\mathcal{S}(\xi)}\mathbf{A}^{2,n}(\xi)$

The TC implies: $\mathbf{A}^{1,n}(\xi) = \rho_1^{DD}(\xi) \mathbf{A}^{2,n-1}(\xi)$

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Convergence analysis (D-D case)

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The TC implies:
$$A^{1,n}(\xi) = \rho_1^{DD}(\xi) A^{2,n-1}(\xi) \longrightarrow A^{1,n}(\xi) = \rho_1^{DD}(\xi) \rho_2^{DD}(\xi) A^{1,n-2}(\xi)$$

$$\mathbf{A}^{2,n}(\xi) = \rho_2^{DD}(\xi) \mathbf{A}^{1,n-1}(\xi) \longrightarrow \mathbf{A}^{2,n}(\xi) = \rho_1^{DD}(\xi) \rho_2^{DD}(\xi) \mathbf{A}^{2,n-2}(\xi)$$

Convergence analysis (D-D case)

$$\begin{aligned} & \mathscr{L}_{H} \mathbf{v}^{1,n} = 0 & \text{in } \widetilde{\Omega}_{1} & \mathscr{L}_{H} \mathbf{v}^{2,n} = 0 & \text{in } \widetilde{\Omega}_{2} \\ & \mathbf{v}^{1,n} = 0 & \text{on } \partial \widetilde{\Omega} & \mathbf{v}^{2,n} = 0 & \text{on } \partial \widetilde{\Omega} \\ & (\partial_{x} + \widetilde{p}_{1,2}) \mathbf{v}^{1,n} = (\partial_{x} + \widetilde{p}_{1,2}) \mathbf{v}^{2,n-1} & \text{on } \widetilde{\Gamma}_{12} & (\partial_{x} + \widetilde{p}_{2,1}) \mathbf{v}^{2,n} = (\partial_{x} + \widetilde{p}_{2,1}) \mathbf{v}^{1,n-1} & \text{on } \widetilde{\Gamma}_{21} \end{aligned}$$

Using separation of variables methods, we get that

$$\mathbf{v}^{i,n} = \sum_{\xi \in \mathbb{N}} N_{\xi} \sin(\xi \pi y) \left(\mathbf{A}^{i,n}(\xi) e^{i\mathcal{S}(\xi)x} + \mathbf{B}^{i,n}(\xi) e^{-i\mathcal{S}(\xi)x} \right) \quad \text{where } \mathcal{S}(\xi) = \sqrt{\widetilde{\mu} - (\xi \pi)^2}.$$

The boundary conditions implies: $\mathbf{B}^{1,n}(\xi) = -\mathbf{A}^{1,n}(\xi)$ and $\mathbf{B}^{2,n}(\xi) = -e^{2i\mathcal{S}(\xi)}\mathbf{A}^{2,n}(\xi)$

The TC implies:
$$A^{1,n}(\xi) = \rho_1^{DD}(\xi) A^{2,n-1}(\xi) \longrightarrow A^{1,n}(\xi) = \rho_1^{DD}(\xi) \rho_2^{DD}(\xi) A^{1,n-2}(\xi)$$

$$\mathbf{A}^{2,n}(\xi) = \rho_2^{DD}(\xi) \mathbf{A}^{1,n-1}(\xi) \longrightarrow \mathbf{A}^{2,n}(\xi) = \rho_1^{DD}(\xi) \rho_2^{DD}(\xi) \mathbf{A}^{2,n-2}(\xi)$$

 $-\rho$ (S)

Convergence analysis (D-D case)



Illustration on an example (Wave Ray: $-\Delta \mathbf{u} + i\mathbf{a} \cdot \nabla \mathbf{u} = 0$, $\widetilde{\omega} = 5$, $\widetilde{p}_{1,2} = \widetilde{p}_{2,1} = i\omega$)



Convergence analysis (D-D case)



Illustration on an example (Wave Ray: $-\Delta \mathbf{u} + i\mathbf{a} \cdot \nabla \mathbf{u} = 0$, $\widetilde{\omega} = 5$, $\widetilde{p}_{1,2} = \widetilde{p}_{2,1} = i\omega$)


Convergence analysis (D-PML case)

$$\begin{aligned} & \mathcal{L}_{H}^{PML} \mathbf{v}^{1,n} = 0 & \text{in } \widetilde{\Omega}_{1} \\ & \mathbf{v}^{1,n} = 0 & \text{on } \partial \widetilde{\Omega} \\ & (\partial_{x} + \widetilde{p}_{1,2}) \mathbf{v}^{1,n} = (\partial_{x} + \widetilde{p}_{1,2}) \mathbf{v}^{2,n-1} & \text{on } \widetilde{\Gamma}_{12} \end{aligned} \qquad \begin{aligned} & \mathcal{L}_{H}^{PML} \mathbf{v}^{2,n} = 0 & \text{in } \widetilde{\Omega}_{2} \\ & \mathbf{v}^{2,n} = 0 & \text{on } \partial \widetilde{\Omega} \\ & (\partial_{x} + \widetilde{p}_{2,1}) \mathbf{v}^{2,n} = (\partial_{x} + \widetilde{p}_{2,1}) \mathbf{v}^{1,n-1} & \text{on } \widetilde{\Gamma}_{21} \end{aligned}$$

Using separation of variables methods, we get that

$$\mathbf{v}^{i,n} = \sum_{\xi \in \mathbb{N}} N_{\xi} \Psi_{\xi}(y) \left(\mathbf{A}^{i,n}(\xi) e^{i\mathcal{S}(\lambda_{\xi})x} + \mathbf{B}^{i,n}(\xi) e^{-i\mathcal{S}(\lambda_{\xi})x} \right)$$

where $(\Psi_{\xi}, \lambda_{\xi})$ are the eigenfunctions and eigenvalues of

$$-\partial_{\widetilde{y}}^{2} \Psi_{\xi} = \lambda_{\xi}^{2} \Psi_{\xi}, \quad y \in [0,1]$$
$$\Psi_{\xi} = 0, \qquad y \in \{0,1\}$$

where \widetilde{y} is the complex stretched coordinate.

Convergence analysis (D-PML case)

$$\begin{aligned} & \mathcal{L}_{H}^{PML} \mathbf{v}^{1,n} = 0 & \text{in } \widetilde{\Omega}_{1} \\ & \mathbf{v}^{1,n} = 0 & \text{on } \partial \widetilde{\Omega} \\ & (\partial_{x} + \widetilde{p}_{1,2}) \mathbf{v}^{1,n} = (\partial_{x} + \widetilde{p}_{1,2}) \mathbf{v}^{2,n-1} & \text{on } \widetilde{\Gamma}_{12} \end{aligned} \qquad \begin{aligned} & \mathcal{L}_{H}^{PML} \mathbf{v}^{2,n} = 0 & \text{in } \widetilde{\Omega}_{2} \\ & \mathbf{v}^{2,n} = 0 & \text{on } \partial \widetilde{\Omega} \\ & (\partial_{x} + \widetilde{p}_{2,1}) \mathbf{v}^{2,n} = (\partial_{x} + \widetilde{p}_{2,1}) \mathbf{v}^{1,n-1} & \text{on } \widetilde{\Gamma}_{21} \end{aligned}$$

Using separation of variables methods, we get that

$$\mathbf{V}^{i,n} = \sum_{\xi \in \mathbb{N}} N_{\xi} \Psi_{\xi}(y) \left(\mathbf{A}^{i,n}(\xi) e^{i\mathcal{S}(\lambda_{\xi})x} + \mathbf{B}^{i,n}(\xi) e^{-i\mathcal{S}(\lambda_{\xi})x} \right)$$

where $(\Psi_{\xi}, \lambda_{\xi})$ are the eigenfunctions and eigenvalues of

$$-\partial_{\widetilde{y}}^{2} \Psi_{\xi} = \lambda_{\xi}^{2} \Psi_{\xi}, \quad y \in [0,1]$$
$$\Psi_{\xi} = 0, \qquad y \in \{0,1\}$$

where \widetilde{y} is the complex stretched coordinate.

Remark: It is not clear that the eigenfunctions $(\Psi_{\xi})_{\xi}$ form a basis of $L^2([0,1])$!

Convergence analysis (D-PML case)

$$\begin{aligned} \mathscr{L}_{H}^{PML} \mathbf{v}^{1,n} &= 0 & \text{in } \widetilde{\Omega}_{1} \\ \mathbf{v}^{1,n} &= 0 & \text{on } \partial \widetilde{\Omega} \\ (\partial_{x} + \widetilde{p}_{1,2}) \mathbf{v}^{1,n} &= (\partial_{x} + \widetilde{p}_{1,2}) \mathbf{v}^{2,n-1} & \text{on } \widetilde{\Gamma}_{12} \end{aligned} \qquad \begin{aligned} \mathscr{L}_{H}^{PML} \mathbf{v}^{2,n} &= 0 & \text{in } \widetilde{\Omega}_{2} \\ \mathbf{v}^{2,n} &= 0 & \text{on } \partial \widetilde{\Omega} \\ (\partial_{x} + \widetilde{p}_{2,1}) \mathbf{v}^{2,n} &= (\partial_{x} + \widetilde{p}_{2,1}) \mathbf{v}^{1,n-1} & \text{on } \widetilde{\Gamma}_{21} \end{aligned}$$

Using separation of variables methods, we get that

$$\mathbf{v}^{i,n} = \sum_{\xi \in \mathbb{N}} N_{\xi} \Psi_{\xi}(y) \left(\mathbf{A}^{i,n}(\xi) e^{i\mathcal{S}(\lambda_{\xi})x} + \mathbf{B}^{i,n}(\xi) e^{-i\mathcal{S}(\lambda_{\xi})x} \right)$$

Similar calculations as before show that

$$\mathbf{A}^{1,n}(\xi) = \rho_1^{DPML}(\xi) \mathbf{A}^{2,n-1}(\xi)$$

$$\mathbf{A}^{2,n}(\xi) = \rho_2^{DPML}(\xi) \mathbf{A}^{1,n-1}(\xi)$$

Convergence analysis (D-PML case)

$$\begin{aligned} \mathcal{L}_{H}^{PML} \mathbf{v}^{1,n} &= 0 & \text{in } \widetilde{\Omega}_{1} \\ \mathbf{v}^{1,n} &= 0 & \text{on } \partial \widetilde{\Omega} \\ (\partial_{x} + \widetilde{p}_{1,2}) \mathbf{v}^{1,n} &= (\partial_{x} + \widetilde{p}_{1,2}) \mathbf{v}^{2,n-1} & \text{on } \widetilde{\Gamma}_{12} \end{aligned} \qquad \begin{aligned} \mathcal{L}_{H}^{PML} \mathbf{v}^{2,n} &= 0 & \text{in } \widetilde{\Omega}_{2} \\ \mathbf{v}^{2,n} &= 0 & \text{on } \partial \widetilde{\Omega} \\ (\partial_{x} + \widetilde{p}_{2,1}) \mathbf{v}^{2,n} &= (\partial_{x} + \widetilde{p}_{2,1}) \mathbf{v}^{1,n-1} & \text{on } \widetilde{\Gamma}_{21} \end{aligned}$$

Using separation of variables methods, we get that

$$\mathbf{v}^{i,n} = \sum_{\xi \in \mathbb{N}} N_{\xi} \Psi_{\xi}(y) \left(\mathbf{A}^{i,n}(\xi) e^{i\mathcal{S}(\lambda_{\xi})x} + \mathbf{B}^{i,n}(\xi) e^{-i\mathcal{S}(\lambda_{\xi})x} \right)$$

Similar calculations as before show that

$$A^{1,n}(\xi) = \rho_1^{DPML}(\xi) A^{2,n-1}(\xi) \longrightarrow A^{1,n}(\xi) = \rho_1^{DPML}(\xi) \rho_2^{DPML}(\xi) A^{1,n-2}(\xi)$$

$$A^{2,n}(\xi) = \rho_2^{DPML}(\xi) A^{1,n-1}(\xi) \longrightarrow A^{2,n}(\xi) = \rho_1^{DPML}(\xi) \rho_2^{DPML}(\xi) A^{2,n-2}(\xi)$$

$$:= \rho^{DPML}(\xi)$$

Convergence analysis (D-PML case)



Illustration on an example (Wave Ray: $-\Delta \mathbf{u} + i\mathbf{a} \cdot \nabla \mathbf{u} = 0$, $\widetilde{\omega} = 5$, $\widetilde{p}_{1,2} = \widetilde{p}_{2,1} = i\omega$)



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Convergence analysis (D-PML case)



Illustration on an example (Wave Ray: $-\Delta \mathbf{u} + i\mathbf{a} \cdot \nabla \mathbf{u} = 0$, $\widetilde{\omega} = 5$, $\widetilde{p}_{1,2} = \widetilde{p}_{2,1} = i\omega$)



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Comparison of the four situations (Wave Ray: $-\Delta \mathbf{u} + i\mathbf{a} \cdot \nabla \mathbf{u} = 0$, $\widetilde{\boldsymbol{\omega}} = 5$, $\sigma_{PML} = 10$)



Comparison of the four situations (Wave Ray: $-\Delta \mathbf{u} + i\mathbf{a} \cdot \nabla \mathbf{u} = 0$, $\widetilde{\boldsymbol{\omega}} = 5$, $\sigma_{PML} = 50$)



Comparison of the four situations (Wave Ray: $-\Delta \mathbf{u} + i\mathbf{a} \cdot \nabla \mathbf{u} = 0$, $\widetilde{\omega} = 5$, $\sigma_{PML} = 50$)



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Comparison of the four situations (Wave Ray: $-\Delta \mathbf{u} + i\mathbf{a} \cdot \nabla \mathbf{u} = 0$, $\widetilde{\omega} = 50$, $\sigma_{PML} = 10$)



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Comparison of the four situations (Wave Ray: $-\Delta \mathbf{u} + i\mathbf{a} \cdot \nabla \mathbf{u} = 0$, $\widetilde{\boldsymbol{\omega}} = 50$, $\sigma_{PML} = 10$)



Comparison of the four situations (Wave Ray: $-\Delta \mathbf{u} + i\mathbf{a} \cdot \nabla \mathbf{u} = 0$, $\widetilde{\boldsymbol{\omega}} = 50$, $\sigma_{PML} = 10$)





1. Motivation

- 2. Link with Helmholtz equation
- 3. Convergence analysis on a toy problem
- 4. An alternated iterative algorithm
- 5. Conclusion

















Convergence in the four situations (Wave Ray: $-\Delta \mathbf{u} + i\mathbf{a} \cdot \nabla \mathbf{u} = 0$, $\widetilde{\boldsymbol{\omega}} = 5$, $\sigma_{PML} = 10$)





Some ideas on the convergence analysis of the alternated algorithm:

 $\mathbf{A}^{i,n}(\xi_k) \longrightarrow \rho(\xi_k) \mathbf{A}^{i,n}(\xi_k) = \mathbf{A}^{i,n+1/2}(\xi_k)$

Step 1

 $\mathbf{A}^{J,n}(\zeta_0)$

Some ideas on the convergence analysis of the alternated algorithm:

 $\mathbf{A}^{i,n}(\xi_k) \longrightarrow \rho(\xi_k) \mathbf{A}^{i,n}(\xi_k) = \mathbf{A}^{i,n+1/2}(\xi_k)$ $\mathbf{A}^{j,n}(\zeta_k)$

Step 1 Ste

Step 2

 $\mathbf{A}^{i,n}(\xi_{k}) \longrightarrow \rho(\xi_{k}) \mathbf{A}^{i,n}(\xi_{k}) = \mathbf{A}^{i,n+1/2}(\xi_{k})$ $\stackrel{i}{:} :$ $\mathbf{A}^{i,n}(\xi_{k}) \longrightarrow \rho'(\xi_{k}) \mathbf{A}^{i,n}(\xi_{k}) = \mathbf{A}^{i,n+1/2}(\xi_{k})$ $\stackrel{i}{:} :$ $\mathbf{A}^{i,n}(\xi_{k}) \longrightarrow \rho'(\xi_{k}) \mathbf{A}^{i,n}(\xi_{k}) = \mathbf{A}^{i,n+1/2}(\xi_{k})$

Step 1 Step 2 Step 3

 $\mathbf{A}^{j,n}(\zeta_0) \longrightarrow \rho'(\zeta_0) \mathbf{A}^{j,n}(\zeta_0) = \mathbf{A}^{j,n+1/2}(\zeta_0)$ $\vdots \qquad \vdots$ Step 2 Step 3 Step 1 Step 4 29



Illustration of the matrix $\mathcal{P}_{12\to AB}$ (Wave Ray: $-\Delta \mathbf{u} + i\mathbf{a} \cdot \nabla \mathbf{u} = 0$)







 $\widetilde{\omega} = 5$

 $\widetilde{\omega} = 50$

 $\widetilde{\omega} = 150$

Illustration of the matrix $\mathcal{P}_{12\to AB}$ (Wave Ray: $-\Delta \mathbf{u} + i\mathbf{a} \cdot \nabla \mathbf{u} = 0$)





Illustration of the matrix $\mathcal{P}_{12\to AB}$ (Wave Ray: $-\Delta \mathbf{u} + i\mathbf{a} \cdot \nabla \mathbf{u} = 0$)



Illustration of the matrix $\mathcal{P}_{12\to AB}$ (Wave Ray: $-\Delta \mathbf{u} + i\mathbf{a} \cdot \nabla \mathbf{u} = 0$)







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1. Motivation

- 2. Link with Helmholtz equation
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Using PML has a strong impact on the convergence of classical iterative algorithm.





Using PML has a strong impact on the convergence of classical iterative algorithm.



We propose an alternated algorithm based on splitting once vertically and once horizontally the domain. This algorithm :

* improve the convergence factor in every case











Using PML has a strong impact on the convergence of classical iterative algorithm.



We propose an alternated algorithm based on splitting once vertically and once horizontally the domain. This algorithm :

- * improve the convergence factor in every case
- * and have a different behaviour depending on the PML BCs,






Using PML has a strong impact on the convergence of classical iterative algorithm.

